Long-term fault slip rates, distributed deformation rates, and forecast of seismicity in the western United States from joint fitting of community geologic, geodetic, and stress-direction datasets

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ABSTRACT. The long-term-average velocity field of the western United States is computed with a kinematic finite-element code. Community datasets include fault traces, geologic offset rates, geodetic velocities, principal stress directions, and Euler poles. There is an irreducible minimum amount of distributed permanent deformation, which accommodates 1/3 of Pacific-North America relative motion in California. Much of this may be due to slip on faults not included in the model. All datasets are fit at a common RMS level of 1.8 datum standard deviations. Experiments with alternate weights, fault sets, and Euler poles define a suite of acceptable community models. In pseudo-prospective tests, fault offset rates are compared to 126 additional published rates not used in the computation: 44% are consistent; another 48% have discrepancies under 1 mm/a, and 8% have larger discrepancies. Updated models are then computed. Novel predictions include: dextral slip at 2–3 mm/a in the Brothers fault zone, two alternative solutions for the Mendocino triple junction, slower slip on some trains of the San Andreas fault than in recent hazard models, and clockwise rotation of some domains in the Eastern California shear zone. Long-term seismicity is computed by assigning each fault and finite element the seismicity parameters (coupled thickness, corner magnitude, and spectral slope) of the most comparable type of plate boundary. This long-term seismicity forecast is retrospectively compared to instrumental seismicity. The western U.S. has been 37% below its long-term-average seismicity during 1977-2008, primarily because of (temporary) reduced activity in the Cascadia subduction zone and San Andreas fault system.

1. Motivation

There are at least two reasons to pursue a unified kinematic model of ongoing deformation in each of the world’s orogens: (1) Dynamic theory and modeling (which involve rheology, stress-equilibrium, and driving forces) will be more nearly correct when they develop from a good kinematic description of what is actually happening. (2) Any complete kinematic model can be converted to a long-term seismicity forecast, from which seismic hazard maps and seismic risk statistics can be computed for guidance of public policy and personal choices.

This paper contributes to both goals. By computing minimum rates of distributed permanent deformation (between model fault traces), I will show that this distributed deformation accommodates a significant fraction of relative plate motion in California, and that kinematic or dynamic models with purely-elastic microplates separated by a small number of plate-boundary faults are not appropriate. By converting the preferred model to a long-term seismicity forecast which is independent of historical seismicity, I highlight regions in which future seismicity will probably be greater than historical seismicity. A subsidiary goal is to
illustrate a process for mapping of long-term seismicity which is rule-based, objective, and transparent, while providing a mechanism for frequent and inexpensive updates as new data become available.

2. Modeling Algorithms, Contrasted with Predecessors

The computational framework for this paper is a set of three codes, each of which has been presented previously with full mathematical detail. Here is a brief qualitative description of each, followed by some distinctions between each program and the methods used by other researchers.

2.1. Program Slippery

The computation of uncertainty in the long-term geologic offset rate from a single offset feature, and also the uncertainty in multi-feature combined offset rates for a particular fault train, is contained in program Slippery.f90 presented by Bird [2007], who included the source code in a digital appendix. (A fault train is a contiguous piece of the trace of a fault system along which our knowledge of fault geometry permits the null hypothesis of uniformity of one component of long-term offset rate.) Each offset distance is classified as one or more of 6 types, depending on the geometry of measurement: R (right-lateral trace-parallel heave), L (left-lateral trace-parallel heave), D (divergent trace-perpendicular heave), P (convergent trace-perpendicular heave), N (normal-sense throw), or T (thrust-sense throw). Oblique offsets are decomposed into two components and treated as two data. The uncertainty in the offset distance measured at the fault trace is represented by a probability density function (PDF) which is typically Gaussian (except in cases of upper and/or lower limits). Uncertainty in the far-field offset is increased by consideration of plausible changes in regional elastic strain, based on amounts of ground-breaking seismic slip which have been observed on other faults of the same type. The age of the offset feature is also represented by a PDF, which may have several different forms depending on whether the age is directly measured or bracketed, and on whether the dating method has problems of inheritance. The PDFs for offset distance and offset age are combined by an integral formula to obtain the PDF for the long-term (far-field) offset rate. From this PDF it is easy to select the median rate (at cumulative probability 0.5), and the lower and upper 95%-confidence limits (at cumulative probabilities of 0.025 and 0.975, respectively). The formal standard deviation is also computed, even though this PDF is not typically Gaussian.

Offset rates from individual offset features can be combined when they lie on the same fault train. First, the program estimates the chance that each individual rate is incorrect, unrepresentative, or inapplicable to neotectonics, using an empirical formula developed in Bird [2007]. Then, the PDFs of individual rates are combined by a formula which considers all weighted combinations of potentially-reliable rates to determine the PDF for the combined offset rate. Again, median rate and 95%-confidence limits are easily obtained from this PDF. The formal standard deviation is also computed, even though this PDF is not typically Gaussian.

While similar calculations involving PDFs have been made by a few authors in studies of single faults, most authors have been content to divide a lower limit on offset at the fault trace by an upper limit on age (and vice versa) to obtain a range of rates for each offset feature. They have rarely considered the complication of plausible elastic strain changes in any systematic way.

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Previous regional seismic hazard studies [e.g., 2007 Working Group on California Earthquake Probabilities, 2008; hereinafter abbreviated as 2007 WGCEP, 2008] have typically decided fault slip rates by deliberation in a committee of experts. While the fastest (and most dangerous) faults received very careful consideration, many slow-moving faults have been assigned uncertainties by rule-of-thumb (e.g., ±25% or ±50% of the selected offset rate), which are almost always too small. Also, these committees have considered additional factors such as kinematic compatibility, plate tectonics, geodetic velocities, paleoseismicity, and historical seismicity when choosing their preferred slip rates. For brevity, I will refer to these as “consensus composite rates.” Consensus composite rates are not appropriate as inputs to NeoKinema (described below), in which these non-geologic factors are also considered and automatically balanced against offset rates which should be purely geologic (even if this leaves them highly uncertain).

2.2. Program NeoKinema

The merger of geologic offset rates, geodetic velocities, and principal stress directions to estimate the long-term velocity field is accomplished with kinematic finite-element code NeoKinema.f90, which was used by Bird & Liu [2007], Liu & Bird [2008], and Rucker [2008]. The equations underlying the program were developed in Supplemental Material S1 (sm001.pdf) of Liu & Bird [2008]. Source code was listed as their Supplemental Material S2 (sm002.zip), but note that this previously-published version (v.2.1, 2007.08.14) is no longer the latest, as described below.

The model domain is the area within a closed curve on the Earth’s spherical surface. The domain is divided into many spherical-triangle finite-elements [Kong & Bird, 1995], with nodes at their corners (Figure 1). The degrees of freedom are two at each node: Southward component of long-term-average velocity, and Eastward component of long-term-average velocity. Therefore, differentiation of velocity within each triangle yields the long-term-average 2-D (horizontal plane) strain rate tensor, which is permanent (not elastic) by definition. The remaining components of the 3-D permanent strain rate tensor are derived from conservation of volume and verticality of one principal axis. It is not necessary to model vertical velocity components explicitly.

The general formalism for solving for nodal horizontal velocity components is to optimize a weighted-least-squares objective function by finding its stationary point in multidimensional velocity-component space with a system of linear equations. Nonlinearities are handled by iteration of the solution (typically 20 times). Velocity boundary conditions are usually applied all around the edges of the models, which should ideally lie within relatively rigid parts of the surrounding plates.

Geodetic benchmarks are treated as internal point constraints on the velocity field (with associated uncertainties). However, geodetic velocities are first “corrected” to remove local elastic bending due to temporary locking of the seismogenic portion of (most) faults, using the current model estimates of the fault slip rates, locking depths assigned a priori, and analytic solutions for rectangular dislocations in a uniform elastic half-space. This requires iteration.

Faults with positive target offset rates contribute to the target strain rates of all elements they cut through. Uncertainty in fault offset rate contributes to anisotropic compliance of all elements that a fault cuts through. An unlimited number of faults can cut through any element, as long as no node lies exactly on a fault trace. However, better accuracy is expected when fast-
slipping faults are outlined by narrow quadrilaterals formed of pairs of elongated triangular elements (Fig. 1). Input fault offset rate components can be either heave rates or throw rates. Throw rates are converted into heave rates using assumed fault dips [Table 5 in Bird & Kagan, 2004]. All dip-slip faults are permitted to slip somewhat obliquely (but restrained by a control parameter) for more realistic flexibility of the fault network. This also requires iteration of the solution.

In elements with no mapped fault traces (“continuum elements”) the horizontal principal directions of the long-term permanent strain rate are constrained by horizontal principal stress directions, which are interpolated from data of the World Stress Map into every finite element by the clustered-data method of Bird & Li [1996]. (Stress-regime information from WSM is not used.) Unfaulted elements also have a target strain-rate of zero, with an assigned uncertainty. This uncertainty [parameter μ of Appendix S1 of Liu & Bird, 2008] is obtained in bootstrap fashion by iteration of the entire solution.

The objective function of NeoKinema is a nondimensional functional of both dimensional model predictions ($p$) and corresponding dimensional data values ($r$), normalized by dimensional covariance matrix ($C$) or by individual datum standard deviations ($σ$):

$$\Pi \equiv -\left(\begin{array}{c} \tilde{\mathbf{p}} - \tilde{\mathbf{r}} \end{array}\right)^T \left[\tilde{\mathbf{C}}^{-1}_{\text{GPS}}\right] \left(\begin{array}{c} \tilde{\mathbf{p}} - \tilde{\mathbf{r}} \end{array}\right) - \frac{1}{L_0} \sum_{m=1}^{M} \int \left(\frac{p_m - r_m}{\sigma_m}\right)^2 \text{d}\ell - \frac{1}{A_0} \sum_{n=1}^{3} \int \left(\frac{p_n - r_n}{\sigma_n^2}\right)^2 \text{d}a$$

where the first term is a quadratic form involving the great vector of all geodetic velocity components and its covariance matrix $\tilde{\mathbf{C}}_{\text{GPS}}$, the second term concerns the $M$ long-term fault offset-rates $r_m$ with their uncertainties $σ_m$, and the third term concerns the constraints on sizes and orientations of distributed permanent deformation-rate tensors (in 2-D, with 3 independent components) in between the mapped faults. Note that this objective function gives a result that is (approximately) independent of the sizes of the finite elements into which the length and area integrals are subdivided.

This objective function includes two “tuning” parameters: (1) trace length for unit weight of long-term offset-rate data, $L_0$; and (2) area receiving unit weight in continuum stiffness and isotropy constraints, $A_0$. (Both are relative to constant unit weight of point-based geodetic data.) Adjustment of these two values controls the relative quality of the fits to geodetic data (best fit with large $L_0$ and large $A_0$), geologic data (best fit with small $L_0$ and large $A_0$), and continuum constraints (including both minimization of strain-rate and orientation of strain-rate; best fit with large $L_0$ and small $A_0$).

The quality of any particular model is described by 3 dimensionless misfit measures, each of which is a root-mean-square norm ($N_2$) of a vector of nondimensionalized misfits to data:

$$N_{2\text{geodetic}}^2 \equiv \frac{1}{2B} \sum_{k=1}^{B} \left(\tilde{\mathbf{p}}_b - \tilde{\mathbf{r}}_b\right)^T \left[\tilde{\mathbf{C}}^{-1}_b\right] \left(\tilde{\mathbf{p}}_b - \tilde{\mathbf{r}}_b\right)$$

where $B$ is the number of geodetic benchmarks and this error measure at each benchmark involves only the local (2×2) covariance of its 2 horizontal components $\tilde{\mathbf{C}}_b$, and
where the $a_i$ are the areas of the finite elements, and the predictions and data are both transformed versions of the azimuth of the most-compressive principal horizontal strain-rate. One important objective in modeling is to bring these measures below ~2, and as close as possible to 1. (Fits with $N_2 < 1$ could be considered overconstrained; there would be some risk of fitting the high-frequency noise in the data as well as its useful low-frequency signals.)

In previous projects we used a parallel measure of the misfit to long-term geologic offset rates, weighted only by trace-lengths (and inversely by datum variances). However, this measure gave potentially misleading results by suggesting a better fit than had actually been achieved. This is due to the very nonuniform populations of fault offset rates. Somewhat like earthquake moments in a seismic catalog, they span many orders of magnitude (e.g., 4.6 orders, from 0.001 mm/a to 40 mm/a, in this project). Also like earthquakes, the small rates are far more numerous than the large rates, which occur on only a few first-order fault trains (San Andreas, Mendocino, Cascadia, etc.). Finally, there is a tendency for many datum standard deviations to be the same order-of-magnitude as the rate (at least for relatively well-constrained rates). A weighted-least-squares algorithm like NeoKinema will always fit those data best which have the smallest standard deviations. So, NeoKinema routinely matches with great precision all of those slow offset rates which also have small standard deviations. An inappropriate misfit measure can make this look like a successful fit to all offset rates, when in fact the fit to the rates of first-order faults may be unacceptable. After some experimentation, I programmed a better misfit measure in which, prior to the $N_2$ (RMS) norm operation, the dimensionless misfits are each weighted by the seismic potency rate of their associated fault. (Seismic potency rate is the product of seismogenic fault area and slip rate.) For stability of this measure, I use the greater of the model or datum slip-rate to determine this relative weight within the misfit measure. This new misfit measure is called the “potency” misfit:

$$N_2^{\text{potency}} \equiv \sqrt{\sum_{i=1}^{\text{elements}} a_i \sum_{m=1}^{M} \ell_{im} w_m h_{im}^{\text{sup}} \left( \frac{p_{im} - r_m}{\sigma_m} \right)^2}$$ (4)
information to constrain the model crustal flow outside of fault zones and increase its dynamic plausibility; this has the practical effect of permitting many small finite elements to be used for better spatial resolution of fault interactions.

The modeling of the western U.S. presented here is most similar to that of Bird & Liu [2007], who used a previous version of NeoKinema. The use of revised misfit measure (equation 4) is the primary change in the algorithm. Other differences in application are that I (1) incorporate faults in the southern Gorda region of the Juan de Fuca plate, and in the Rio Grande rift; (2) use new geologic and geodetic compilations with reliable uncertainties; and (3) perform more tests of model sensitivity to Euler poles, fault sets, weighting factors, and new data. These differences will each be developed in following sections of this paper.

2.3. Program Long_Term_Seismicity

Program Long_Term_Seismicity.f90 is a realization of the set of hypotheses known as the SHIFT model (an acronym for Seismic Hazard Inferred From Tectonics) [Bird & Liu, 2007]. The primary hypotheses are that: (1) The long-term seismic moment rate of any tectonic fault, or any large volume of permanently-deforming lithosphere, is approximately that computed using the coupled seismogenic thickness of the most comparable type of plate boundary. (2) The long-term seismicity of any tectonic fault, or any large volume of permanently-deforming lithosphere, is approximately that computed from its moment rate using the frequency-magnitude distribution of the most comparable type of plate boundary. The seismicity coefficients (coupled seismogenic lithosphere thickness \(cz\), corner magnitude \(m_c\), and asymptotic spectral slope \(\beta\) of the tapered Gutenberg-Richter frequency/moment relation) of each type of plate boundary were determined by Bird & Kagan [2004] and listed in their Table 5. Decision rules for assigning faults and finite elements to the “most comparable” type of plate boundary are contained in Tables 1 & 2 of Bird & Liu [2007].

Recent analysis of global seismicity by Bird et al. [2009?] has shown that the earthquake productivity of subduction zones and continental convergent boundaries is nonlinear in relative plate velocity. This revision is incorporated in version 3 of Long_Term_Seismicity, which was used in this project.

The primary difference between this method and that of recent seismic hazard forecasts for California [e.g., 2007 WGCEP, 2008] and the western U.S. [e.g., Frankel et al., 1996, 2002; Petersen et al., 2008] is that I never assume that faults have either periodic or characteristic earthquakes, and I do not assume that earthquake magnitude is limited by mapped fault length or inferred fault area. Instead, I propose that (with low probability) an earthquake beginning on a short fault, or in an area between mapped faults, can grow to large size by linking up mapped faults and/or existing-but-unmapped faults, and occasionally by creating new fault area [Black, 2008]. The practical result of this difference in assumptions can be seen by comparing the RELM seismicity forecasts mapped by Field [2007], especially his Figures 3.1 and 3.2 compared to 3.9. Another difference is that my method does not use historical seismicity or inferred paleoseismicity of the region in any direct way. Recent seismicity is an important consideration in short-term forecasting, but I consider that seismic catalogs (whether historic or instrumental) are too short, and paleoseismic catalogs presently too incomplete, to provide a sound basis for long-term seismicity projections.
2.4. Availability of Codes

Source code for program Slippery was in Bird [2007]. Fortran 90 source codes for NeoKinema (v.2.2, 2008.01.30) and Long_Term_Seismicity (v.3, 2009.04.29) are attached to this publication as supplemental materials: NeoKinema_v2p2_Guadalupe.f90.txt and Long_Term_Seismicity_v3.f90.txt. All source codes used in this project are also available from the author at: http://peterbird.name, where there are also accessory programs, including OrbWin for creation of 2-D spherical F-E grids, OrbNumber for renumbering nodes to reduce bandwidth, NeoKineMap for graphical display of input and output datasets, and RangeFinder for summarizing the fault offset rates predicted in a suite of successful NeoKinema models.

3. Community Datasets and Other Inputs

Most of the calculations presented in this paper are based on datasets created by others in long-standing collaborative groups, including the Working Group[s] on California Earthquake Probabilities, Southern California Earthquake Center, USGS National Seismic Hazard teams, Plate Boundary Observatory geodesists, and World Stress Map team. Therefore, they are referred to here as “community models” (although I retain responsibility for any errors in assumptions or computation).

3.1. Traces of active and potentially-active faults

Traces of active and potentially-active faults in the western U.S. and adjacent offshore regions were compiled from 5 sources:

Fault traces in California (and its continental borderland) are from Fault Model 2.1 or 2.2 of the Working Group on California Earthquake Probabilities (Figure 2). As explained in 2007 WGCEP [2008], these resulted from the merger of (1) the Community Fault Model [Plesch et al., 2007] created by the Southern California Earthquake Center, with (2) traces in northern California adopted or created by WGCEP [2003]. Fault Models 2.1 and 2.2 are mutually exclusive alternatives which differ primarily in the shapes and topologies of certain fault traces in the southern margin of the Transverse Ranges, from the Santa Barbara Channel eastward to the Puente Hills of California. They have 243 and 248 traces, respectively. A community Internet application named SCEC-VDO (Southern California Earthquake Center-Virtual Display of Objects) may be used to display these faults in 3-D. The Fault Models contain estimated locking depth ranges, which in southern California are largely from Nazareth & Hauksson [2004]. (Consensus composite slip rates are also included in the Fault Models, but were not used in this project.) NeoKinema fault numbers (e.g., “F4170”, used in Table 4 and in the supplemental files attached to this paper) were assigned by adding 4000 to WGCEP fault numbers. Two faults which are common to both Fault Models have internally inconsistent data which make it unclear whether they were intended to be oblique-slip thrusts or purely strike-slip faults: the San Andreas (San Gorgonio Pass-Garnet Hill) train has dip of 58°NE and rake of 180°, while the Santa Rosa Island fault has dip of 90° and rake of 30°. In each case, I covered both possibilities by making the fault purely strike-slip in one model, and treating it as an oblique thrust in the other model.

Fault traces in other western states include all those used in computations for the 2002 National Seismic Hazard Maps [Haller et al., 2002].
I included additional potentially-active faults outside California from my own compilation of the geologic literature [Table 1 of Bird, 2007], including faults with known Neogene activity which lack documented overlap formations. This was based on the consideration that active faults of modest slip rate (e.g., 0.1 mm/a) and typical slip-per-event (e.g., 4 m) may have experienced last movement in the late Pleistocene (e.g., 40 ka), but their scarps may have been obscured by later Pleistocene erosion and/or sedimentation. Many of these faults were identified by authors of regional survey papers about the Basin and Range province or the Rio Grande rift [e.g., Stewart, 1978, 1998; Tweto, 1979; dePolo, 1998], while others were identified during dissertation or other mapping projects reported in the literature. I digitized these additional traces from various sources including state geologic maps, online maps of the USGS Quaternary Fault and Fold Database, and large-scale maps in dissertations and journals. Where a normal fault has a mapped surface trace in Quaternary deposits along only part of a basin/range topographic scarp, I typically assumed that an underlying fault extends along the entire scarp. Likewise, I often combined groups of minor faults into a single “fault system” trace, appropriate for small-scale modeling, where the gaps are small enough to be jumped by earthquake ruptures [Wesnousky, 2006; Black, 2008]. Faults of less than 10 km length which could not be integrated with other nearby traces into a longer fault system were not included.

Traces of the Cascadia subduction zone and the spreading centers and transform faults along the Gorda Ridge are from the PB2002 plate boundary model of Bird [2003].

The 545 fault traces within the Gorda orogen part of the Juan de Fuca plate are from Chaytor et al. [2004], who mapped them using high-resolution swath bathymetry and seismic reflection profiling. These are a combination of reactivated normal faults originally created at the Gorda Rise, and newer faults which cross-cut the seafloor-spreading fabric. Faults of ambiguous slip were assumed to be left-lateral.

All of these 1479 traces (Figure 2) are contained in file fGCN_merged_WGCEPFM2p2_200810.dig.txt which is part of the supplemental material for this paper. The NeoKinema convention is that fault traces are digitized left-to-right when looking in the downdip direction; vertical strike-slip faults are mostly digitized from W to E.

3.2. Long-term geologic offset rates on faults

NeoKinema requires a prior (input) offset rate and uncertainty for each component of slip on each modeled fault. At the end of the computation, it provides a posterior (output) offset rate for each component of slip on each modeled fault. For brevity, the prior (input) rates will also be referred to as “target” rates, and the posterior (output) rates will be referred to as “predicted.”

One distinguishing feature of this model is that it uses no consensus composite slip rates for faults on land, but only geologic offset rates based on dated offset features. The computation of the probability density function (PDF) for the combined long-term offset rate of any fault train with program Slippery.f90 was described briefly in section 2.1, and fully in Bird [2007].

The target offset rates and uncertainties for NeoKinema are the median rate and the formal standard deviation, respectively, from the combined-rate lines of Tables 1 and 2 of Bird [2007]. Rates for California fault trains come from Table 2, which was based on the PaleoSites database addition to the USGS Quaternary Fault and Fold Database, created through the efforts of the Working Group on California Earthquake Probabilities. While this database is not yet
available on-line, it has been reviewed by 3 WGCEP members, as well as by the author and a coworker. Rates for faults in other western states come from Table 1 of Bird [2007], which was based on the author’s personal compilation from the literature. This was reviewed only during the publication process, and the chances of errors and omissions are correspondingly higher.

The total number of geologic offset rates is 572, while the number of fault trains in the model is 1479. Fortunately, NeoKinema is able to model faults that have very uncertain target rates, and to predict their rates from the merger of geodetic, plate-tectonic, stress-orientation, and strain-compatibility considerations. In order for this to work properly, the faults with no documented offset features should be assigned large uncertainties in offset rate, with some rational basis. Such faults are here assigned a generic rate PDF, median rate, and (large) standard deviation based on the composite PDF for all faults of that type (R, L, N, D, T, or P) in the western U.S. which do have dated offset features. For example, a normal fault (N) with no offset datum is assigned a target throw rate of $N = 0.183 \text{ mm/a}$ with a standard deviation of $0.343 \text{ mm/a}$. A right-lateral strike-slip fault (R) with no offset datum is assigned a target heave rate of $R = 6.18 \text{ mm/a}$ with a standard deviation of $12.6 \text{ mm/a}$. These large uncertainties permit the fault to slip much faster or slower than the nominal rate, to remain locked, or even to slip in the opposite sense from the target rate.

In most parts of the NeoKinema calculation it is not important whether a fault slips seismically or aseismically. However, this makes a difference when correcting geodetic velocities of benchmarks near a fault for temporary fault locking, as no correction is needed for faults which creep steadily. In the input data file, certain California faults are designated as creeping by a logical flag: Calaveras (Central, South), Concord, Green Valley (North, South), Hayward (North, South), Hunting Creek-Berryessa, Maacama-Garberville, and San Andreas (creeping segment). It is not known whether other faults outside California might also be creeping, but the distinction is less important when the heave rate of the fault is comparable to or less than the uncertainty in GPS velocity.

Target rates and uncertainties for spreading segments (offsets of type D) and adjacent transforms (offsets of type L, R) on the Gorda Ridge are from magnetic anomaly bands, according to the data compilation of DeMets et al. [1990], corrected for the magnetic timescale revision of DeMets et al. [1994], and interpolated where necessary using latitude as the independent variable. The Cascadia subduction zone (the only offset of type S) is assigned a nominal rate of $39.5 \text{ mm/a}$ [Bird, 2003] with a standard deviation of $7.5 \text{ mm/a}$ to allow for unknown amounts of deformation in the overriding lithosphere.

Faults in the oceanic lithosphere of the Gorda orogen [Chaytor et al., 2004] are probably a distinct population from continental normal and strike-slip faults, with different distribution(s) of rates. Unfortunately, only a single offset has been identified: a long sinistral fault has known minimum slip rate of $(1.5\sim1.7 \text{ km})/(<2 \text{ Ma})$, implying a minimum rate of $(0.75\sim0.85) \text{ mm/a}$. There is not enough information to employ program Slippery. Rather arbitrarily, I reduced the multiple activity classes of Chaytor et al. to only two: “active fault” with target rate of $0 \text{ mm/a}$ and standard deviation of $1 \text{ mm/a}$, and “potentially-active fault” with target rate of $0 \text{ mm/a}$ and standard deviation of $0.3 \text{ mm/a}$. Given this great uncertainty, it would be very valuable to obtain a few seafloor velocities by geodetic means [Chadwell & Spiess, 2008] within the Gorda orogen.

All 1536 offset rates with their uncertainties are compiled in file fGCN_merged_WGCEPFM2p2_200810.nki.txt which is part of the supplemental material for this
paper. The number of rate entries is greater than the number of fault trains because some fault trains are known to have oblique slip, which is described by a strike-slip entry (offset type L or R) plus a separate dip-slip entry (N or D for extension; alternatively T or P for shortening).

### 3.3. Interseismic velocities of benchmarks from GPS

Velocities of benchmarks in California are from a new combined solution of GPS data completed by Zhengkang Shen, Bob King, Min Wang, and Duncan Agnew in June 2006 for the Working Group on California Earthquake Probabilities. A preliminary (November 2005) version of this solution is available from the WGCEP site at: [http://wgcep.org/](http://wgcep.org/). It is a statewide solution based on analysis of the original data (SCEC and Berkeley reprocessed regional survey mode daily solutions, and SOPAC processing of the continuous sites), rather than adjustments of other investigators' velocity fields. Coseismic effects of the Joshua Tree, Landers, Northridge, Hector Mine, and San Simeon earthquakes have been estimated and excluded from the velocity modeling. Data showing immediate short-term (a few months to a year or so) postseismic deformation were also excluded. This solution includes 1226 benchmarks and a covariance matrix.

To provide coverage of other western states, I used the Plate Boundary Observatory joint GPS solution of 2007.09.19 from [http://pboweb.unavco.org/](http://pboweb.unavco.org/). This is the network velocity field derived from final combined solutions generated by the Analysis Center Coordinator at MIT. Only individual-site uncertainty ellipses are available for this solution. I selected sites from this model in four steps: (1) deletion of stations with velocity standard deviations exceeding 3 mm/a (which eliminates most stations with short occupation history and/or nonlinear movement history); (2) deletion of all benchmarks in Yellowstone National Park, which may be affected by magma chamber deflation; (3) deletion of 3 benchmarks (P075 in NV, P683 in ID, P692 in OR) which have anomalous velocities suggesting possible fault-creep or non-tectonic processes; and (4) deletion of all benchmarks in California, which is already covered by the WGCEP solution described above. This left 307 benchmarks, 193 of which are within the domain of the NeoKinema model.

This composite GPS velocity field of 1419 benchmarks is plotted in Figure 3. Both component solutions are expressed in the reference frame of stable (eastern) North America. Certainly there must have been small procedural differences in the definitions of this reference frame by the two groups of geodesists, and this could result in artificial velocity shear across the inland borders of California. However, no discontinuities are apparent (except across active faults of the Walker Lane), and it is likely that any such discrepancy is less than 1 mm/a.

Before using this velocity field in NeoKinema, all benchmarks located less than 2 km from faults with slip rates over 1 mm/a were deleted, because at smaller distances F-E grid GCN8p9.feg interpolates and smears the fault discontinuities in long-term velocity, making it erroneous to compare grid velocity with corrected (long-term) geodetic velocity. This editing step removed 212 or 209 benchmarks, depending on whether WGCEP Fault Model 2.1 or 2.2 was used. Thus, 1207 or 1210 benchmarks were actually used in each NeoKinema solution.

### 3.4. Most-compressive horizontal stress azimuths

For the study of Bird & Liu [2007], most-compressive horizontal principal stress directions were downloaded from the World Stress Map Project [Reinecker et al., 2004; Heidbach et al., 2008]. About 963 data fell inside the Gorda-California-Nevada orogen, and an
additional 1105 data were outside its margins but close enough (<22° of great-circle arc) to be used for the interpolation of principal stress directions. The same dataset is used here. The NeoKinema input file s_Gorda-Cal-Nev.nki.txt is part of the supplemental information attached to this paper.

The uncertainties reported by WSM for each azimuth are highly generalized and somewhat arbitrary. Each datum has a letter-coded quality class, and approximate angular uncertainties are stated for each quality class for each type of data. However, it is unclear whether these numerical values are standard deviations, 95%-confidence limits, or absolute limits. Also, the rounding of these values suggests that they may be subjective estimates rather than results of statistical studies. Therefore, the uncertainties from WSM were not used. Instead, NeoKinema interpolates stress direction to the center of each finite element, using the clustered-data algorithm of Bird & Li [1996] which provides individual uncertainties for each result which are based on the scatter in surrounding azimuths. Standard deviations range from 2.7° to 49.4°, with median of 8.5°. Both original and interpolated stress directions are shown in Figure 4.

3.5. Boundary conditions

Velocity boundary conditions may be imposed around the margins of a NeoKinema simulation, and this is highly desirable as a way of enforcing both (approximate) rigidity of the surrounding plates and correct net relative velocity across the model domain. Because F-E grid GCN8p9.feg (Figure 1) spans the entire Gorda-California-Nevada orogen and Rio Grande rift, it is surrounded by relatively rigid portions of the North America (NA), Pacific (PA), and Juan de Fuca (JF) plates. I take stable NA as the velocity reference frame. Then, neotectonic Euler poles for the northeastern margin of PA, and for JF, are needed in relation to stable NA.

The neotectonic Euler pole for NA-PA is uncertain and controversial, as shown in Figure 5. (All of these NA-PA Euler poles are detailed in Table 2.) Apparent disagreements may reflect Pacific plate deformation and reference-frame issues as well as measurement error. In a global dynamic finite-element model, Bird et al. [2008] computed internal Pacific strain-rates of order 10^-18/s caused by regional stress (mainly NW-SE tension) acting on the olivine-dominated rheology of oceanic lithosphere. These strain-rates integrate to internal relative velocities of only ~0.3 mm/a; however, this olivine rheology has not been tested outside the lab and could be too strong. Probably more significant is thermal contraction, especially in the younger eastern portions of PA; Kumar & Gordon [2009] estimate that this causes relative velocities of 3~10 mm/a. Because internal deformation may not be negligible, I concentrate on finding the Euler pole that will best approximate the motion of the northeastern margin of PA, where it abuts this model.

The NUVEL-1A pole of DeMets et al. [1994] came from a global solution for the poles of the 12 largest plates, based on marine magnetic anomalies 2A, transform fault azimuths, and seismic slip vectors. If PA is internally deforming, this pole should best describe the motion of its eastern parts along the East Pacific Rise. The other poles shown are geodetic, and unfortunately none of them took the NUVEL data set into consideration.

The quality of a purely-geodetic pole depends upon: (1) length of observation; (2) technical issues concerning reference frame and data reduction; and (3) number and locations of sites which represent each plate. Length-of-observation is obviously better for the poles.
published more recently. On the other hand, Argus [2007] raised an important criticism of
conventional plate-motion solutions based on ITRF2000 or (especially) ITRF2005 because these
reference frames drift with respect to the global shell of lithosphere. When a poleward-drifting
reference frame is used to extract horizontal velocity components for Euler-pole calculations, an
equatorial belt of anomalous velocity is introduced which will contaminate NA-PA poles in
particular. This concern was recently addressed by Kogan & Steblkov [2008] with their “plate-
frame” pole. An additional consideration is that the only geodetic pole to represent PA by
benchmarks on Guadalupe Island and Baja California (Figure 3) is that of Gonzalez-Garcia et al.
[2003]. Thus, this geodetic pole could be the best to represent the motion of the northeastern
margin of the Pacific plate, even if the Kogan & Steblkov pole is a more accurate representation of
the motion of the central parts of the Pacific plate.

The JF-PA Euler pole used (35°N, 26°E, 0.5068 degree/m.y.) is from Wilson [1988]. It is
based directly on magnetic anomaly bands along the Juan de Fuca Ridge, and is relatively
certain. The Sierra Nevada/Great Valley plate of Argus & Gordon [2001] is entirely included
within the model domain (Figure 1) and does not require boundary conditions.

3.6. Fixed parameters

Certain additional parameters read as input by NeoKinema were fixed throughout this
modeling project. Each solution was iterated 20 times, which typically resulted in overall
relative velocity changes of 0.03% and maximum velocity changes of ~0.1 mm/a in the last
iteration. Parameter \( \xi \) (\( \xi \)), a small strain-rate quantum used in the code to prevent
singularities, was fixed at 3.2×10^{-17} /s based on previous experience. The standard deviation of
slip rakes for dip-slip faults (around the nominal target of \( \pm 90^\circ \)) was 20°. The shallower and
dereeper interseismic locking depth limits for faults outside California were 1 and 12 km,
respectively, except in the Cascadia subduction zone where they were 14 and 40 km [Bird &
Kagan, 2004]. All optional program features were switched off.

4. High Rates of Distributed Permanent Deformation

Distributed deformation is defined here as that part of the field of strain-rate tensors in a
NeoKinema solution which is not due to slip on modeled faults. It is permanent by definition
because NeoKinema solves for long-term-average (10^4- to 10^6-year) velocities, and cyclical
variations in elastic strain average to insignificant rates over many earthquake cycles. In the
weighted-least-squares algorithm of NeoKinema, distributed deformation rate is treated as an
undesirable error and minimized. However, solutions with real data show a recalcitrant residual
which cannot be eliminated. For plotting, tabulation, and discussion it is convenient to convert
strain-rate tensors to scalars; for consistency with the objective function of NeoKinema, I use the
azimuthally-invariant scalar measure:

\[
\dot{\varepsilon} = \sqrt{\dot{\varepsilon}_{NS}^2 + \dot{\varepsilon}_{NS} \dot{\varepsilon}_{EW} + \dot{\varepsilon}_{EW}^2 + \dot{\varepsilon}_{NE}^2}
\]

which in strike-slip regimes (\( \dot{\varepsilon}_{NS} + \dot{\varepsilon}_{EW} = 0 \)) is equal to the greatest horizontal principal strain
rate, or the shear strain rate in fault-parallel coordinates. Spatial variations of scalar \( \dot{\varepsilon} \) are
mapped in Figure 6. It can be further characterized by its area-weighted RMS value, \( \mu^* \).

Because NeoKinema only computes a model of the surface velocity field, it is unclear
how deep this distributed deformation extends. However, the current paradigm for continental
tectonics is that the seismogenic layer is bounded by a brittle/ductile transition, below which
distributed deformation by climb-assisted steady-state dislocation creep is widespread. Thus, we
may say that the “problem” or “innovation” of distributed permanent deformation is primarily
notable in the seismogenic layer, extending to perhaps 12 km depth. In the following
subsections I will discuss why shallow distributed deformation must exist, what strain
mechanisms might be involved, how its intensity is constrained by this study, and what this
implies for future kinematic and dynamic modeling.

4.1. Arguments and observations concerning distributed deformation

If all crustal blocks were completely outlined by faults, including transform faults
conforming to arcs of small circles about Euler poles, there would still be some distributed
deformation in the regions surrounding unstable triple-junctions [McKenzie & Morgan, 1969].
Actual distributed deformation in continents must be greater because many faults simply end
where their slip goes to zero. In the Basin and Range province it is often possible to estimate the
(minimum) throw on normal faults from the height of their topographic scarps, and it is common
to see that throw taper to zero at fault ends, which do not connect to transform faults as plate
theory predicts. Some strike-slip faults (which are mapped in the plane of motion) also end
without connections; examples in WGCEP Fault Model 2.1 include the Hosgri (Extension),
Ortigalita, Greenville (South), Santa Ynez, Ludlow, Earthquake Valley, and Owens Valley
faults. (This list does not include cases of aligned but widely-separated faults where a cryptic
connection is possible.)

Another kind of evidence for distributed deformation is the well-known discrepancy
between (higher) geodetic rates of dextral strain and (lower) geologic rates of dextral strain based
on measured fault slip rates in the Eastern California shear zone [Oskin et al., 2007]. (In this
case time-dependence of regional deformation has been proposed as an alternative explanation,
but no physical mechanism for time-dependence has been modeled.)

There is also a theoretical dynamical argument for distributed deformation, based on the
rheologic layering of the lithosphere: A brittle/ductile transition at midcrustal depth can only be
maintained if the ductile layer has a non-zero permanent strain rate to give it strength; without
distributed deformation the ductile layer would undergo viscoelastic relaxation, transferring
deviatoric stress upward into a thinning brittle layer until this would eventually break [Roy &
Royden, 2000]. Thus, even a plate with a shallow frictional layer cannot sustain deviatoric stress
for a million years if its regional strain-rate is zero.

A few mesoscale investigations have identified shallow distributed deformation in
favorable circumstances, especially in close proximity to major faults. Jamison [1991] mapped
en-echelon folds developed in transpression along the San Andreas, Rinconada, and Newport-
Inglewood faults; later, Argus & Gordon [2001] used his results to infer at least 0.8±0.5 mm/a
dextral and 1.1±0.6 mm/a compressional deformation adjacent to the San Andreas. Salyards et
al. [1992] used paleomagnetic declinations to estimate distributed deformation in marsh deposits
around the San Andreas fault at Pallet Creek; off-fault deformation was 3 times greater than fault
slip. (However, this is usually considered a special circumstance specific to marshy
paleoseismic sites.) Unruh & Lettis [1998] studied seismogenic transpressional deformation east
of the Hayward fault, where local fold-and-thrust belts are proposed to shorten at several mm/a.
Flodin & Aydin [2004] studied distributed deformation of the Aztec Sandstone by strike-slip
faults in Valley of Fire State Park, Nevada where they identified 5 hierarchical generations of
structures in outcrop. Oskin et al. [2007] used the inactive Silver Bell normal fault as a strain marker in the belt surrounding the active Calico dextral fault, and found distributed deformation within 500 m of the Calico trace which was ~23% of total offset. One possibility is that distributed deformation is rare, and these sites have been described because they are exceptional. Another is that distributed deformation is common, but is not typically so visible in outcrops. This would depend on whether distributed deformation is typically accomplished by faulting, or by true continuum deformation.

4.2. Strain mechanisms of distributed deformation

At least 5 different strain mechanisms might contribute to distributed permanent deformation at shallow depths. (1) Silicate crystals accommodate small amounts of strain by primary transient creep (“cold work”), which is nonequilibrium dislocation glide ending in tangles unrelieved by climb [Poirier, 1985]. This kind of strain occurs at a declining rate following the first imposition of deviatoric stress, and due to “work hardening” it is much less in subsequent loading cycles. (2) Rocks whose dominant minerals are stable below the water table (quartz, calcite) can deform by solution transfer, as is commonly seen in folded sedimentary rocks [Gratz, 1991]. (3) Rocks whose dominant minerals are not stable (mafic rocks) can deform by an analogous but non-steady weathering process, in which stressed grains and grain corners are preferentially weathered to create unstressed clays. (4) Tensile microcracking which is due to differential expansion of different minerals (and differently-oriented crystals of the same mineral) during erosional unloading [Bruner, 1984] can have a preferred orientation where there is also a regional deviatoric stress of tectonic origin [Boness & Zoback, 2006]. (5) Distributed deformation can occur by frictional sliding on many faults of small net offset which have not been included in the NeoKinema model.

A critical distinction is that mechanisms (1)-(4) would be practically aseismic, while mechanism (5) could produce damaging earthquakes (although at low rates). The best available double-difference relocations of California earthquakes [Hauksson & Shearer, 2005] continue to show that many earthquakes cannot be located on any of the faults of the Community Fault Model. Therefore, it is prudent to base seismic hazard estimates on the hypothesis that slip on unmodeled faults is the dominant mode of distributed deformation. There is an analogy between quasi-fractal fault networks and the power-law distribution of earthquake moments, in both of which the majority of strain is due to first-order features, but all scales make some contribution. The faults modeled in this paper are simplified from primary observations recorded on geologic maps, and quadrangle-scale geologic maps rarely represent more than a fraction of the faults actually present in the field. In this light, the quantity $\mu^*$ can be considered an artifact of our limited progress in mapping and modeling, rather than a fundamental physical property of the crust. (However, it is likely to be a practical reality in modeling for centuries to come.)

4.3. Constraining minimum $\mu^*$ with NeoKinema

NeoKinema models are best computed in sets, because it is necessary to find the optimal values for 3 critical input parameters: $L_0$, $A_0$, and $\mu$. The first two should be adjusted to find acceptable fits to geologic, geodetic, and stress-direction data simultaneously, as described in section 2.2. Meanwhile, an input/prior value of $\mu$ (the model parameter) must be found which will produce a similar output/posterior value of $\mu^*$ (the computed RMS rate of scalar distributed...
deformation rate $\dot{e}$ in that model). Fortunately, experience shows that there is a neighborhood
around the optimum point in 3-D parameter space where $\mu^*$ is relatively insensitive to these 3
inputs.

The reconnaissance models described here were performed without using the full
covariance matrix of the geodetic velocities in California; only the block-diagonal part
(individual site error ellipses) was used. This reduced run times from 36 hours to 75 minutes
each. Fault Model 2.1 and the NUVEL-1A pole for NA-PA were used throughout this set.

Having previously determined that $\mu^* \approx 5 \times 10^{-16}$ s by trial-and-error, I ran a systematic
set of 45 models in which $\mu = 5 \times 10^{-16}$ /s was fixed, while $L_0$ was increased from 1250 to
320000 m by factor-of-2 steps, and $A_0$ was increased $2 \times 10^8$ to $32 \times 10^8$ m$^2$ by factor-of-2 steps
(Table 3). Inside this rectangle in 2-D parameter space, an elliptical region was found (Figure
7) in which 10 models had acceptable misfit measures of $N^2_{\text{geodetic}} < 2$, $N^2_{\text{potency}} < 2$, and
$N^2_{\text{stress}} < 2$ simultaneously. Figure 8 shows resulting $\mu^*$ values with contours: there is a flat
region with minimum $\mu^* > 5 \times 10^{-16}$ /s in the lower right, and 4 acceptable models have $\mu^*$
below $6 \times 10^{-16}$ /s.

Slicing in the orthogonal direction through parameter space, Figure 9 and Table 3 show
results of 8 more models in which weights ($L_0 = 4 \times 10^4$ m, $A_0 = 4 \times 10^8$ m$^2$) were fixed while
prior/input $\mu$ was varied. First, it is clear that posterior/output $\mu^*$ is only weakly dependent on
input $\mu$, so that $5 \times 10^{-16}$ /s is the only value that gives consistency of prior with posterior.
Second, we see that even if we abandon consistency and try to force less distributed deformation,
all 3 misfit scores quickly rise to unacceptable values because of the increased rigidity of the
model. For this geographic region, with these input data, there are no successful models with
$\mu^*$ below $5 \times 10^{-16}$ /s.

4.4. Implications for kinematic and dynamic modeling

I used the preferred model from this project (GCN2008088, described below) to create a
budget for right-lateral deformation along the San Andreas plate boundary. The PA-NA
transform plate boundary stretches 1350 km in this model, from the Mendocino triple junction
(124.41°W, 40.26°N) to the northwestern Gulf of California (114.21°W, 31.35°N). According to
the Guadalupe pole for NA-PA (Table 2), the mean velocity along this boundary is 47.8 km/m.y..

The product of these numbers is 64530 (km)$^2$/m.y.. However, the line integral of dextral slip
rates on dextral and dextral-oblique (offset type R) faults in the model, to the southeast of the
Mendocino triple junction, is only 43235 (km)$^2$/m.y., or 67% of this. I also computed the area-integral of twice the dextral strike-slip distributed permanent strain-rate on vertical planes
trending N38°W, also to the SE of the Mendocino triple-junction: the result was 21940
(km)$^2$/m.y., or 34% of the total. Thus, slip on mapped faults takes up 2/3 of PA-NA relative
motion in the latitudes of the San Andreas fault system, while distributed deformation takes 1/3.
(Clockwise rotation does not seem to play a significant part when averaged across the state of
California, although it is locally important as discussed below.)

This conclusion conflicts with that of Humphreys & Weldon [1994], who summed
geologic slip rates on 3 paths across southern California to equal total PA-NA relative velocity.
The paths they selected may not be typical. Also, their study did not incorporate geodetic data, so it was possible to attribute missing deformation to offshore fault systems which lack geologic constraints. Shen-Tu et al. [1999] also computed a model of the western United States that was based on 100 “geologic slip rates” and matched PA-NA relative velocity. However, 44 of their rates were from Petersen & Wesnousky [1994], who supplemented missing geologic slip rates with consensus composite rates that often reflect geodesy, seismicity, and/or kinematic compatibility (assuming rigid microplates). Another 14 of their rates from other sources had similar non-geologic bases. Neither of these studies included independent statistical analysis of geologic offsets and their ages comparable to that in Bird [2007].

Interpretations of geodetic velocities have traditionally assumed purely-elastic microplates in the seismogenic depth range (although they vary according to the rheologies and layering assumed at greater depth). Since this assumed shallow structure can be described by 2 parameters per fault (slip rate and locking depth) it is relatively easy to determine both by inversion. Now I propose that many fault-crossing profiles are also sampling significant amounts of either distributed permanent deformation (if it is aseismic and steady in time) or elastic straining preparatory to future distributed permanent deformation (if it is seismic and unsteady in time). Since NeoKinema models (e.g., Figure 6) predict that distributed deformation is often concentrated near major faults, it may be quite difficult to distinguish between these models using geodetic velocities alone. One prediction of this new model is that, on average, inversions of geodetic velocities using elastic microplates have tended to overestimate locking depths; this can be determined in subsequent earthquakes, although postseismic deep creep is a confusing factor. Another approach is to continue collecting and refining geologic offset rates, to see if traditional inversions of geodetic data have tended to overstate the slip rates of the dominant faults, as I expect. Perhaps it will even be possible to directly invert for the fraction of distributed deformation where its spatial distribution can be independently constrained. For this purpose, histograms of well-located seismicity (averaged along a fault, and plotted in cross-section) might serve as a reasonable proxy.

This finding also has important implications for dynamic modeling. If fault systems are quasi-fractal, and only ~2/3 of long-term-average deformation is accommodated on those master faults included in community models, then dynamic models which use purely-elastic microplates cannot be expected to succeed. It will be necessary to use crustal blocks which are plastic or frictional, and to make the rheologies of fault zones and inter-fault blocks as similar as possible. Dynamic finite element programs like Shells [Bird, 1999] take this approach, keeping dislocation-creep rheology laterally uniform, and merely distinguishing the fault from the rest of the crust by a lower effective coefficient of friction. Much work will be needed to understand why and how effective friction drops as crust is progressively deformed (or whether some other kind of unified rheology is more appropriate).

5. Experiments with Euler Poles, Fault Models, and Geodetic Covariance

Details of the 53 models described in Section 4.3 (Figures 7~9) are presented in Table 3. They all used the NUVEL-1A Euler pole (approximating relative rotation between stable NA and the northeastern margin of PA) and WGCEP Fault Model 2.1 in California. The best model in the group was GCN2008028, with misfits of: \( N^\text{geodetic}_2 = 1.730 \), \( N^\text{potency}_2 = 1.523 \), and \( N^\text{stress}_2 = 1.746 \); its overall misfit level is rated as: \( \sup (N^\text{geodetic}_2, N^\text{potency}_2, N^\text{stress}_2) = 1.746. \)
The NUVEL-1A pole gives the highest azimuth for relative velocity of stable NA with respect to stable PA when computed at Parkfield, CA (Table 2); that is, it gives the most transpressional model. In models GCN2008-053~057 and -094, I investigated the results of recomputing the boundary conditions using each of 5 published geodetic poles: ITRF2000, PA_GPS, Guadalupe, ITRF2005, and Plate_frame. The comparison of NUVEL-1A with Plate_frame is especially interesting because they span the range from most transpressional to most transtensional model; the range of azimuths for relative plate velocity is 4°. They are also the “slowest” and “fastest” of the poles modeled in terms of their predictions at Parkfield: the range is 45.7~50.4 mm/a.

By differencing output files from these two extreme models and plotting differential velocity, I found that the changes are mostly offshore. This is natural, as velocities on land are strongly constrained by GPS velocities which are relative to stable eastern NA. Therefore, the Plate_frame pole gave about 4.7 mm/a of additional dextral shear along the length of the borderland, and about 3.3 mm/a more compression perpendicular to the borderland. We have very little offshore data that is relevant to choosing between these models. However, there were slight variations in on-land velocities in 4 regions: Point Arena, San Francisco peninsula, Salinia terrane, and the southern tip of the model domain in northern Baja California. The Plate_frame pole causes the San Andreas (North Coast) dextral rate to increase from 14.7 to 17.5 mm/a, coming closer to its geologic target rate. The San Andreas (Peninsula) dextral rate increases from 14.5 to 17.3 mm/a, increasing its overrate error. The San Andreas (Santa Cruz Mt.) dextral rate increases from 17.7 to 19.6 mm/a, but is still well below its geologic target. The dextral rate on the Rinconada fault increases from 0.05 to 2.8 mm/a. The San Gregorio (North & South) dextral rates also increase by 1.5~2 mm/a. All other rate changes on faults for which we have geologic constraints are less than 1 mm/a.

The effect of the Euler pole on overall model misfit is very modest, and occurs by changing $N_2^{stress}$ which is uniformly the highest misfit in this set of 6 models. The best result (Table 3) is not for either NUVEL-1A or Plate_frame but for the Guadalupe pole which lies between them (Figure 5). As mentioned above, this pole is also the only geodetic pole to incorporate Pacific-plate sites closer to California than Hawaii. For both reasons, the Guadalupe pole of Gonzalez-Garcia et al. [2003] was selected as the best for representing the northeastern margin of PA in this modeling project, and employed to fix boundary velocities in most of the subsequent calculations.

The next experiment was to keep all input parameters constant, and retain the Guadalupe pole, but to use WGCEP Fault Model 2.2 in southern California (GCN2008056 vs. -055 in Table 3). Right slip on the San Andreas (San Gorgonio Pass-Garnet Hill) train dropped from 11.1 to 5 mm/a when its throw rate was reduced from 0.6 to 0 mm/a by treating it as a purely strike-slip fault (see section 3.1 above). Right slip on the Mission Creek train increased from -0.05 to 2.2 mm/a in partial compensation. Right slip on the adjacent San Andreas (Coachella) train decreased from 16.6 to 15.2 mm/a. Right slip on the Brawley seismic zone dropped from 12.8 to 8.2 mm/a due to its modified shape. Right slip on the Imperial fault dropped from 24.6 to 18.6 mm/a, presumably due to the modified Brawley seismic zone. Left slip on the Malibu Coast fault increased from 2.1 to 3.2 mm/a due to its new trace. Left-transpressional faulting in the Santa Barbara Channel region was reorganized, and implausible extensional slip on two thrust faults in Fault Model 2.1 was eliminated. Most other changes in slip rate were either under 1 mm/a, or occurred on faults which are only present in one of the Fault Models. The effect on
misfit measures was mixed, with $N_{\text{stress}}^2$ going down slightly and the other two measures rising slightly.

The next innovation was to use the full covariance matrix of the California GPS velocities, which increases NeoKinema run times greatly. Using Fault Model 2.2 and the Guadalupe pole, I computed 18 models with varying $L_0$ and $A_0$, while fixing $\mu$ at $5 \times 10^{-16}$ /s. Results in Figure 10 and Table 3 show the same pattern as in Figure 7, only slightly offset in parameter space. Seven of these models were successful according to the criterion,

$$\sup \left( N_{\text{geodetic}}^2, N_{\text{potency}}^2, N_{\text{stress}}^2 \right) < 2.$$

Then another 5 models (GCN2008077–081, Table 3) tested all other combinations of the two Fault Models with the NUVEL-1A, Guadalupe, and Plate_frame poles. Neither model with the Plate_frame pole succeeded, because their geodetic misfits $N_{\text{geodetic}}^2$ rose to 2.1–2.2. Lastly, 6 more models (GCN2008082–087, Table 3) repeated all the models that had succeeded using the Guadalupe pole and Fault Model 2.2, but now using Fault Model 2.1.

This concludes the set of models which I refer to as the “community” models because of the origins of their input data. By testing various combinations of Fault Models, Euler poles, and NeoKinema weighting parameters, a suite of 16 “acceptable community models” has been found (counting only the latter computations using the full geodetic covariance within California). This provides a good estimate of the ranges of fault offset rates which might be obtained while still fitting all of the community input data reasonably well.

6. Pseudo-prospective Tests, and Updated Models

A true prospective test of these models will require collection of new geologic, geodetic, and/or stress-direction data following publication of this paper. However, there is an opportunity to conduct a pseudo-prospective test immediately by examining the prediction of data already published but not used in the computation.

By surveying major journals through October 2008, I located 54 additional papers giving 126 “new” offset rates on 68 fault trains in the western U.S. (beyond those tabulated by Bird [2007] and used in the community models). All are detailed in Table 4. The column “Community model predictions” gives the range of long-term offset rates predicted by the set of 16 acceptable community models just described. The columns “New geologic offset rate: Min., Max.” give the 95%-confidence limits on long-term fault offset rate obtained by analyzing the new offset in program Slippery.f90 (as described in section 2.1 above) as an individual feature. New rates are plotted against predictions in Figure 11.

The best results among these 126 pseudo-prospective tests were the 44% which showed no discrepancy. In these cases there was some overlap between the 95%-confidence bounds on the new offset rate (as computed by Slippery.f90 considering only the single offset feature) and the range of predictions among the 16 acceptable community models. For example, 2 new provisional dextral rates on the San Andreas (Mojave S) train of 5.9–18.5 and 11–57 mm/a from Weldon et al. [2008] overlap the 16.2–17.4 mm/a range of model predictions. On the same fault train, the new rate of 16–29.5 mm/a based on alluvial fan #3 of Matmon et al. [2005] has no discrepancy (but their rates based on other fans do, as described below). On the San Andreas (San Bernardino N) train one provisional new rate of 7.2–20 mm/a from McGill et al. [2008]
overlaps the model range of 18.9–20.6 mm/a. On the adjacent San Andreas (San Bernardino S) train, another provisional new rate from McGill et al. [2008] of 8.1–21.7 mm/a includes the model range of 11.6–15.4 mm/a. On the San Andreas (Coachella) train, the new dextral rate of 12–16.4 mm/a from Behr et al. [2008] overlaps the community model predictions of 14.8–17.5 mm/a. Off the San Andreas system, there were some cases where the NeoKinema community models predicted the offset rates of faults even in the absence of any offset geologic features to constrain their rates: right slip on the Owens Valley dextral/normal fault was predicted to be 1.44–2.18 mm/a, and is actually estimated as 0.63–5.3 mm/a [Lee et al., 2001b]. Convergent heave on the Compton blind thrust was predicted as 0.85–1.56 mm/a, and is actually estimated as 1.3–2.7 mm/a [Dooling et al., 2008]. (This is an incomplete list; see Table 4 for other cases.) Another 48% of these pseudo-prospective tests resulted in “small” discrepancies of less than 1 mm/a. Another new provisional dextral rate of 13–18 mm/a on the San Andreas (San Bernardino N) train [McGill et al., 2008] is slightly discrepant with the model prediction range of 18.9–20.6 mm/a. Admittedly, there are many cases where the same discrepancy would be “large” if stated as a percentage. For example, the model predictions of 0.18–0.23 for normal throw rate on the Carson Range normal fault miss the new geologic rate of 0.88–15 mm/a [Ramelli et al., 1999] by a discrepancy of only 0.65 mm/a, but by a large fraction. On the other hand, within this subset of 60 small discrepancies, there are 33 cases (55%) in which the community models were not guided by any dated offset geologic features in the tables of Bird [2007] that provided their geologic targets. Predicting a fault offset rate in advance of any geologic measurement is a hard test for any deformation model. The remaining cases are 11 (8%) in which the discrepancy was larger than 1 mm/a. Nine of these were under 4 mm/a, and two were much larger (11 and 26 mm/a, respectively). These problems cluster in the Mojave Desert region of California, where it is well-known that geologic and geodetic rates are difficult to reconcile. These large discrepancies will be considered individually in the regional discussion below.

After this comparison, the 126 new offset rates (and 10 new fault traces) were combined with those already published [Tables 1 & 2 of Bird, 2007] to obtain updated combined geologic target rates and standard deviations for all fault trains. Four additional NeoKinema models, here called the “updated” model set, were computed with the revised geologic target rates and uncertainties and fault traces, keeping other input datasets unchanged. These models used either the NUVEL-1A or the Guadalupe pole for NA-PA, and either WGCEP Fault Model 2.1 or 2.2 in California. On the basis of minimum misfit, model GCN2008088 (Guadalupe pole, Fault Model 2.2) was selected as the best updated model, and therefore as the “preferred” model of this paper, which is displayed in most map-view figures. Figure 12 shows the long-term velocity field of this model.

It is interesting how little the preferred model changes as a result of these 68 updated geologic target rates. Comparing models GCN2008060 and -088, no offset rate changes by more than 3 mm/a. On the San Andreas fault, the Big Bend train slows from 15.4 to 13.6 mm/a, the Mojave N train speeds up from 17.4 to 20 mm/a, the San Bernardino N train slows from 18.9 to 16.6 mm/a, the San Bernardino S train speeds up from 12.2 to 13.4 mm/a, and the Coachella train speeds up from 15.1 to 17 mm/a. On the Elsinore fault, the Temecula stepover train speeds up from 0.8 to 3.7 mm/a, and the en-echelon Glen Ivy stepover train slows from 3.7 to 1.3 mm/a in local compensation. Throw rate increases on the Carson Range normal fault from 0.2 to 2.9 mm/a. All other changes are less than 1.8 mm/a. (All updated offset rates for fault trains with
new data are shown in Table 4.) In this test, the addition of several years of new (or newlycatalogued) geologic rates had only modest effects on the preferred NeoKinema model, indicating its stability. On the other hand, this stability means that many discrepancies remain: 10 (8% of new data) remain above 1 mm/a, and 25 (20% of new data) below 1 mm/a. This is desirable behavior if the discrepancies are due to errors [Bird, 2007] in the new data, but not desirable if the errors are in the model.

7. Regional Discussion and Ad-hoc Experiments

NeoKinema provides predictions of fault offset rates in two formats. In most text and in Table 4 of this paper, the quantities described as model predictions have been the length-weighted along-trace averages of model offset rates in all the finite elements cut by one fault train. When this along-trace average is plotted all along the trace, as in Figure 13A, the result is a ribbon of uniform width; this is easy to interpret, but potentially misleading. The more informative display of Figure 13B shows per-element estimates of fault offset rate, with all their noisy discontinuities in strike and value. Such a display includes some artifacts (especially unreasonably high rates at some fault terminations), but also displays some important variations in slip-rate along traces which are due to interactions between faults and/or distributed permanent deformation. In the remainder of this paper, detailed/noisy plots similar to Figure 12B will be shown in order to convey more information.

7.1. Washington and Oregon

As in most plate-tectonic models, relative motion in the greater Washington-Oregon region (including adjacent seafloor) is dominated by spreading/transform activity on the northern Gorda Ridge and convergence in the Cascadia subduction zone (Figure 13). (Spreading on the Juan de Fuca Ridge is not shown in this figure because it is outside the model domain; see Figure 1). Convergence in the Cascadia subduction zone is relatively constant at values near the mean rate of 36.7 mm/a (Figure 13A, B). However, the strike-slip component changes locally with the azimuth of the trace of the plate boundary, so that its mean of 2.8 mm/a dextral (Figure 13A) conceals local variations from 15 mm/a sinistral to 29 mm/a dextral (Figure 13B). (These local variations have little tectonic significance; I mention them only to illustrate the difference between these two methods of plotting the predictions of the same model.)

In the Cascadia forearc offshore Oregon (43°–46.5°N) 11 WNW-trending high-angle faults have been mapped by Goldfinger et al. [1992] and/or Personius et al. [2003]. All were entered in my database as nominally sinistral faults, although this was based primarily on the conceptual model of Goldfinger et al.; only the Wecoma and Coos Basin faults have sinistrally-offset features, and only the Wecoma fault has a geologic offset rate, of 9.1±2.2 mm/a, which comes from the part of the fault on the Juan de Fuca plate. In the preferred NeoKinema model, sinistral motion on the Wecoma fault is preserved because of its relatively well-constrained geologic rate, but most of the other faults are predicted to slip with a dextral sense, at lower rates. This raises serious doubt about the continuity of the Wecoma fault where it crosses the Cascadia trench. The part of the fault on the Juan de Fuca plate is sinistral where it offsets the Astoria fan, but perhaps the part of the fault in the North America plate is dextral, and the alignment between these two opposite-sense faults is coincidental and temporary (as the Juan de Fuca plate drifts NE relative to NA). The kinematic incompatibility that would normally arise between aligned sinistral and dextral faults would be relieved in this case by a quadruple-junction with the
Cascadia trench, and a decrease in subduction rate on the N side of this junction relative to the S side.

An interesting prediction of the preferred NeoKinema model is that Oregon is bisected by an active dextral fault system composed of 5 aligned faults: from NW to SE, the Tillimook Bay fault (predicted mean dextral rate 1.9 mm/a), the Newburg fault (3.8 mm/a), the Mount Angel fault (3.2 mm/a), the Clackamas River fault (3.2 mm/a), the Sisters fault zone (0.5 mm/a), and the 280-km-long Brothers fault zone [Lawrence, 1976; Walker, 1977; Christiansen & Yeats, 1992] (3.0 mm/a). Distributed deformation bridges the gaps between these traces to create a continuous belt of dextral shear at about 3 mm/a. In the model, this belt acts as a strike-slip transfer (tear) fault system accommodating the northern termination of many normal faults in southeastern Oregon [Lawrence, 1976] or northwestern Nevada (Figures 12, 13). Because none of these predicted dextral faults had a well-constrained geologic slip rate in the input data, this result is primarily dictated by regional kinematic compatibility, and by the PBO GPS velocity solution. Additional campaign-mode GPS velocities that can check this prediction (because they were not included in the “community” datasets) were published by Hammond & Thatcher [2005]. They interpreted clockwise relative rotation between their microplate CSOR (Central Southern OR) and stable NA with Euler pole (-118°E, 44.3°N, -0.8°/m.y.) that would be consistent with dextral slip on the Brothers fault zone, at rates increasing from 2.1 mm/a at its SE end to ~3 mm/a at its NW end (where there would also be an extensional component). However, Hammond & Thatcher did not discuss this fault system, or identify any other discrete microplate boundary. To test this part of the NeoKinema model, I computed one ad-hoc model (GCN2008100) with the addition of 49 new Hammond & Thatcher [2005] GPS velocities in Oregon to those used previously. This model scored slightly better than the “preferred” model GCN2008088 because of its lower stress misfit (which was probably due to the “dilution” of the influence of a questionable high GPS velocity at MDMT in the WGCEP solution). Model GCN2008100 predicts a mean dextral slip rate of 2.3 mm/a instead of 3 mm/a on the Brothers fault, but in every other way is qualitatively identical to the preferred model. This is another demonstration of the stability of the NeoKinema modeling process.

Another area of relatively rapid faulting in this region is the thrust belt in the seaways of Juan de Fuca Strait, San Juan archipelago, and Puget Sound. The West Coast-West San Juan-Survey Mountain thrust along the SW side of Vancouver Island has predicted mean heave rate of 0.9 mm/a, with slip beginning at Clayoquot and increasing southeasterward to 1.3 mm/a in southeastern Vancouver Island. To the east, this shortening is divided between a north branch on the Devils Mountain thrust (mean heave rate 0.44 mm/a; locally up to 0.65), and a south branch on the South Whidbey Island thrust (mean heave rate 0.63 mm/a). Further south in Puget Sound, a crustal block between the Seattle and Tacoma thrust faults is predicted to be elevated at throw rate 0.2 mm/a by heave rates of about 0.55 mm/a on each of these faults. (Other faults not mentioned have lower mean rates.) This association of active thrusting with deep glacial troughs is intriguing. Perhaps it is due to an observer bias resulting from higher population densities and/or easier access in these areas. Or, if it is real, it could reflect an enhancement of thrusting by the Pleistocene glacial removal of topographic mass that would otherwise oppose and moderate thrusting. A similar process on a grander scale was proposed for the Chugach-Wrangell Mountains region of Alaska by Bird [1996].
7.2. Mendocino triple junction region

Relative motion between the rigid northern part of the Juan de Fuca plate and the Pacific (JF-PA) is parallel to the Blanco fracture zone at azimuths of 110~120°, as Chadwell & Spiess [2008] recently confirmed with seafloor geodesy. The Mendocino fault is part of the same JF-PA plate boundary, but has azimuth 93°. This creates a problem of excess crustal area in the southern “Gorda orogen” part of the Juan de Fuca plate.

One possibility is that the Mendocino fault is an oblique right-transpressional fault, with underthrusting of Juan de Fuca crust to the south [Silver, 1971], especially in the Gorda Escarpment portion east of 126°W. One possible indicator of thrusting is a linear dipolar gravity anomaly of 90 mGal amplitude along the Mendocino fault [Leitner et al., 1998] with more negative anomalies to the N and more positive to the S. Another is depression of the Moho and crustal thickening to 12 km in the northeast corner of the Pacific plate within 25 km of the Mendocino triple-junction [Henstock & Levander, 2003] without accompanying surface deformation. Because of these arguments, I permitted oblique slip on the Mendocino fault in most models; in the preferred model, its mean convergent heave rate is predicted to be 10 mm/a, superimposed on a mean dextral rate of 33.5 mm/a (Figure 14A). Therefore, the predicted azimuths of slip vectors would be about 110°, and this is kinematically close to a rigid-plate solution. The many faults mapped by Chaytor et al. [2004] are active in sinistral and/or reverse senses in this model, but mostly at very slow rates of less than 0.1 mm/a. The average rate of the 24 “active” sinistral faults is 0.12 mm/a.

Other authors [Smith et al., 1993; Gulick et al., 2001] have denied any component of thrusting on the Mendocino fault. Because the question is open, I also computed ad-hoc model GCN2008101 in which the Mendocino fault is treated as a purely strike-slip vertical fault. Results in Figure 14B are subtly different: 10 mm/a of N-S shortening is absorbed about equally by fault slip and distributed deformation within the southern “Gorda orogen” part of the Juan de Fuca plate. The average slip velocity on the 24 active sinistral faults identified by Chaytor et al. [2004] increases to 0.28 mm/a. This ad-hoc model also involves a slightly greater indentation of the northeast corner of the Pacific plate, and a slightly reduced slip rates on the northernmost (Offshore) train of the San Andreas fault (7.8 instead of 9.3 mm/a at Cape Mendocino; mean 8.3 instead of 8.8 mm/a). However, there is no dramatic change predicted that would be easy to test on land.

There is another space problem in the region. As pointed out by McCrory [2000], the northernmost part of the San Andreas trace does not align with the Mendocino triple junction, but instead lies ~70 km East of its ideal position. Where the San Andreas bends sharply westward in the King Range/Punta Gorda area, a corner of the Pacific plate is colliding with Cascadia forearc of the North America plate. The SW-dipping King Range and Petrolia thrust faults at this critical corner may not have moved since Early Quaternary time [Jennings, 1994], and are not included in the WGCEP Fault Models. However, extending for 100 km North along the California coast is an active fold-and-thrust belt of mostly NE-dipping thrusts (and blind thrusts beneath anticlines) whose offset rates (or structural growth rates) were catalogued by McCrory [2000]. She estimated their total shortening rate conservatively as 10 mm/a. (My alternative analysis, assuming thrust fault dips of only 20°, suggests that shortening is permitted to range from 14~24 mm/a.) Either is less than PA-NA relative velocity of ~48 mm/a; but the facts that collision is oblique and that many of these thrust faults are longer than the 70-km width
of the indentor may allow for area-balancing. The preferred model of Figure 14A satisfies 9 of McCrory’s 17 rates (Table 4), with a mixture of under- and over-predictions in the other 8 cases. The Russ thrust fault has the only large discrepancy, with predicted throw rate of 3.1–3.5 mm/a exceeding one of McCrory’s two constraints, while nearly agreeing with the other. The model also predicts its second-greatest concentration of distributed deformation (second only to the Imperial Valley region) around Cape Mendocino (Figure 6).

7.3. San Francisco Bay area and central California Coast Ranges

Model predictions (Figure 15) in this area are for dextral slip unless otherwise noted. At 39°N (e.g., Point Arena) the 34.1 mm/a of shear between the borderland and the Great Valley plate is divided among: San Andreas (North Coast) 12.6, Maacama-Garberville 10.2, Bartlett Springs 7.8, and distributed deformation of 3.5 mm/a. At 38°N (e.g., Point Reyes) the 35.3 mm/a of shear is divided among: San Andreas (North Coast) 23.0, Hayward (No) 6.7, Concord 3.0, and distributed deformation of 2.6 mm/a. At 37°N (e.g., Santa Cruz) the 36.6 mm/a of shear is divided among: San Gregorio (No) 0.8, Zayante-Vergeles 1.4, San Andreas (Santa Cruz Mtn) 21.8, Calaveras (So) 4.9, Ortigalita 3.0, and distributed deformation of 4.7 mm/a. At 36°N (e.g., Kettleman City) the 36.6 mm/a of shear is divided among: Hosgri 1.6, Rinconada 0.9, San Andreas (Parkfield) 31.5, and distributed deformation of 2.6 mm/a. (This last exceeds the minimum distributed dextral deformation of 0.9±0.5 mm/a which Argus & Gordon [2001] inferred from mapping of Jamison [1991] in the Temblor Range. The total off-San Andreas dextral shear of 5.1 mm/a at 36°N agrees with the 5±4 mm/a discrepancy of Argus & Gordon.) Note that distributed deformation is only 7–13% of total in this region because of the generally subparallel and continuous fault traces which are also nearly parallel to relative plate motion.

Thrusting is predicted at low convergent heave rates on 4 faults in the region: Mount Diablo thrust 0.17 mm/a, Monte Vista-Shannon thrust 0.26 mm/a, Zayante-Vergeles thrust 0.32 mm/a, and Monterey Bay-Tularcitos thrust 0.9 mm/a (combined with 0.6 mm/a dextral slip).

Comparing mean slip rates in this preferred model with those selected by 2007 WGCEP [2008] for their seismic hazard forecast, the biggest contrast is that this model tends to have lower mean rates on many (but not all) trains of the San Andreas system. From NW to SE, predictions of this model vs. WGCEP include: Offshore train 8.8 vs. 24 mm/a, North Coast 16.2 vs. 24, Peninsula 17.9 vs. 17, Santa Cruz Mtn 22.6 vs. 17, Creeping Segment 29.1, Parkfield 31 vs. 34, Cholame 26.4 vs. 34, Carrizo 25 vs. 34, and Big Bend train 13.6 vs. 34 mm/a. This is because WGCEP rates were primarily based on rigid microplate models, whereas this model has large fractions of PA-NA relative motion accomodated by distributed deformation (see section 4.4 above).

7.4. Southern California

Predicted fault heave rates from the preferred model are shown in Figure 16. Neotectonics in southern California are complicated by the 154-km left step of the San Andreas fault system, which requires thrust-faulting. One organizing factor is the boundary condition applied to the base of the crust by the symmetrical downwelling of mantle lithosphere under the Transverse Ranges [Bird & Rosenstock, 1984]. However, in the absence of true subduction, locations of thrusting change through geologic time due to relative advection of faults, growth of topographic resistance, and growth of bending-stress resistance. Another chaotic or
disorganizing factor is the frequent reactivation of diverse faults formed in earlier stages of the
tectonic history [Ingersoll & Rumelhart, 1999].

A budget for the total rate of thrust-faulting in the Transverse Ranges (from the Tehachipi
Mountains on the N to the San Joaquin Hills on the S) is obtained by multiplying the width of
this left step by the relative velocity of the Pacific plate with respect to the Sierra Nevada/Great
Valley plate: 154 km × 35 km/m.y. = 5390 (km)^2/m.y.. The following 10 thrust faults, listed
with their lengths and mean convergent heave rates, are the most prominent contributors to area
loss among the 75 nominal thrust or oblique-thrust faults in the Transverse Ranges, and together
they make up 50% of the budget: Red Mountain 100 km × 7.1 mm/a = 13.3%, White Wolf 64
km × 6.6 mm/a = 7.8%, Oak Ridge (Offshore) 38 km × 6.7 mm/a = 4.7%, Oak Ridge (Onshore)
49 km × 5.1 mm/a = 4.7%, Santa Susana (alt 2) 43 km × 5.3 mm/a = 4.2%, Simi-Santa Rosa 39
km × 4.6 mm/a = 3.3%, Santa Cruz Island 69 km × 2.6 mm/a = 3.3%, White Wolf (Extension)
46 km × 3 mm/a = 2.6%, Channel Islands thrust 59 km × 2.3 mm/a = 2.6%, and San Cayetano 42
km × 3.1 mm/a = 2.4%. Distributed deformation takes up 38.4% of the budget. The other
11.6% is the net shortening among the other 65 nominal thrust faults in the Transverse Ranges,
but as some of these are predicted to have extensional slip in the preferred model (e.g., Mission
Ridge-Arroyo Parida-Santa Ana, Nacimiento, and San Gabriel) there is some cancellation of area
changes within this group.

It is interesting that very little of the shortening is taken up along the impressive
mountain fronts of the San Bernardino and San Gabriel Mountains (North Frontal faults 77 km
×1.7 mm/a = 2.3%; Mission Creek 32 km × 1.2 mm/a = 0.7%; Cucamonga fault 28 km × 2.6
mm/a = 1.4%); instead it occurs primarily within the lower topography of the Santa Barbara
Channel, Channel Islands, Simi Valley, and San Gabriel Valley. This may be a sign of very low
crustal strength, and the consequent regulation of thrusting by the topographic resistance that it
eventually generates.

Argus et al. [2005] constrained anthropogenic motions with SAR in order to better
analyze GPS velocities in the Los Angeles area, and identified a 25-km-wide belt south of the
San Gabriel Mountains front in which there is 4.5±1 mm/a of crustal shortening. Although they
inferred the Puente Hills thrust to be the most active, in this model the Puente Hills thrust fault
system is well-constrained by geologic data (Table 4) and absorbs only P = 1.4 mm/a of this
shortening. Other active thrusts in this area include the Santa Monica (alt 2) sinistral thrust (L =
1.3, P = 1.3 mm/a), the Hollywood sinistral thrust (L = 1.9, P = 0.8 mm/a), the Raymond sinistral
thrust (L = 2, P = 2.8 mm/a), the Upper Elysian Park thrust (P = 0.8 mm/a), the Lower Elysian
Park dextral thrust (R = 1.3, P = 2.0 mm/a), and the Compton thrust (P = 1.8 mm/a). (These
rates should not be added, as most named faults are shorter than the width of the area described.)
The seismic hazard from thrust faulting is similar to that estimated by Argus et al., but it appears
to be more widely distributed across this urban area.

In the Imperial Valley region of southeastern California and northern Baja California, the
preferred model predicts an maximum dextral shear rate of 38 mm/a. This is less than the 45±2
mm/a that Fialko [2006] inferred from InSAR data (constrained by GPS and EDM data).
Perhaps the radar line-of-sight range rates were affected by long-wavelength vertical movements
of natural or industrial origin. (The discrepancy is only 2 mm/a in range rate.) Alternatively, the
NeoKinema model may predict insufficient fault slip (and too much distributed deformation)
because it uses an incomplete set of fault traces. It is notorious that the primary fault trains of the
plate boundary in this region (Cerro Prieta, Imperial, San Andreas, San Jacinto (Superstition
Mountain), Laguna Salada, and Elsinore (Coyote Mountain) are not mapped as connecting to each other. One reason is tillage. Another is recurring coverage by lacustrine or marine sediments. A third may be tectonic decollement on weak evaporite horizons within the sedimentary section. A fourth is rapid intrusion of basaltic dikes (analogous to seafloor spreading) beneath the sedimentary cover, which may link some mapped faults by creating gaps in the lithosphere. The WGCEP Fault Models attempted to close one large gap between traces by elevating the intrusive center known as the “Brawley seismic zone” to the status of a fault (for strike-slip only), but they left other gaps. Consequently, the preferred model has extremely high rates of distributed deformation in this region (Figure 6). It is important to include this distributed deformation as a potential source of seismicity, which would otherwise be underestimated. If heat-flow or seismic tomography should show the lithosphere to be very thin in some areas, this can be considered when converting rates of distributed permanent deformation to long-term seismic moment rates.

7.5. Mojave Desert: San Andreas fault versus Eastern California shear zone

There has been extensive debate about the long-term crustal flow in this region, which can be summarized by reference to 5 conceptual models. The primary contenders are: (1) a “geologic model” based on dated offset features in the Eastern California shear zone which indicate a low rate of dextral shear [e.g., 5.9±1.4 mm/a per Oskin et al., 2006]. If the motion of the western Mojave relative to stable NA is only 6 mm/a, then the slip rate of the Mojave trains of the San Andreas must be high [e.g., 30±10 mm/a per Matmon et al., 2005; 30–46 mm/a per Rust, 2005]. This model is opposed by a (2) “geodetic model” which estimates Eastern California shear zone motion as 12±2 mm/a [e.g., Sauber et al., 1994] and estimates a lower rate of dextral slip on the Mojave trains of the San Andreas [e.g., 14.3±1.2 mm/a per Meade & Hager, 2005]. Both of these models are conceived in terms of the steady flow of elastic microplates.

Three more concepts attempt to add degrees of freedom to resolve the controversy. Concept (3) “distributed deformation” might reconcile the geodetic rate with geologic offsets in the Eastern California shear zone [Oskin et al., 2007]. Concept (4) the “cyclic model” suggests that crustal flow switches between two modes, and that geologists and geodesists have observed different parts of the cycle [e.g., Dolan et al., 2007]. Concept (5) the “rheologic model” attempts to show how high rates of slip on dextral faults could be disguised by rheologic structures at depth to appear as low rates in simplistic inversions of geodetic data [Dixon et al., 2003; Johnson et al., 2007].

Program NeoKinema is not able to test concepts (4) or (5) because they conflict with key assumptions underlying the program. (Mode-switching has no articulated cause or mechanism, and its kinematics outside the Mojave region are vague. The Dixon et al. [2003] model of the Eastern California shear zone and the Johnson et al. [2007] model of the San Andreas require that there are no faults at the level of the mantle lithosphere, which is especially hard to reconcile with 240 km of net displacement on the southern San Andreas [Buesch & Ehlig, 1982].) So, NeoKinema predictions are necessarily some mixture of models (1), (2), and (3). Results of the preferred model in this study are closest to the “geodetic model” (2) with an added component of “distributed deformation” (3). I do not think that this is due to inadequate weight on the geologic constraints. (See section 2.2 for discussion of how the geologic misfit measure was redefined upward, and section 5 for discussion of how geologic and geodetic misfits were balanced.)
Instead, it is because of two basic constraints: (i) Geologic slip rates are not uniformly high for the Mojave South train of the San Andreas. Including all sources (Ehlert & Ehlig [1977], Buesch & Ehlig [1982], Sieh [1984], Barrows et al. [1985], Frizzell et al. [1986], Meisling & Weldon [1986], Schwartz & Weldon [1986], Sieh et al. [1989], Salyards et al. [1992], Weldon et al. [2002, 2004, 2008], and Matmon et al. [2005]) in an analysis by program Slippery, the combined rate is only 21.9 ± 3.85 mm/a (median ± standard deviation; 95%-confidence range 16.2–29.3 mm/a). (ii) Geodesy has convincingly demonstrated that the Sierra Nevada/Great Valley plate moves NW at 12–13 mm/a relative to stable North America [e.g., Argus & Gordon, 2001]. Because the western Mojave overthrusts the Sierra Nevada/Great Valley plate on the left-transpressional White Wolf and White Wolf (Extension) faults, its velocity to the NW must be higher than this.

The preferred model GCN2008088 has mean slip rates on the Mojave N and Mojave S trains of the San Andreas of 20.1 and 17.4 mm/a, respectively (Figure 16). It is important to note the apparent conflict with recent geologic rates by Matmon et al. [2005] and Rust [2005] which had 95%-confidence lower limits of 21 mm/a (fan #0), 43 (fan #1), 16 (fan #3), 21 (fan #4), 28 (fan #5) and 30 mm/a, respectively. However, each of these rates is only as good as its offset distance, and each offset distance is only as good as the assumption that the drainage crossed the fault in a straight line at a right angle during deposition of the dated sediment. If sediments were deposited at a time when the drainage already had a right-lateral kink, then offset distances and rates have been overestimated. Similar (but left-lateral) arguments may apply to the Garlock (Central) sinistral fault, where the model predicts only 3.8 mm/a, but two offsets identified and dated by McGill & Sieh [1993] imply minimum rates of 5 or 6.2 mm/a, respectively.

Other “large” (> 1 mm/a) discrepancies (Table 4) occur on the Blackwater fault in the Eastern California shear zone, where the model predicts a dextral slip rate of 1.8 mm/a which is higher than two geologic rates of Oskin & Iriondo [2004] with upper limits of 0.3 and 0.5 mm/a, respectively. Since their offset lava flows are pre-Quaternary (7.2 and 3.8 Ma, respectively), a resolution may be possible if Blackwater fault slip began about 1 Ma.

Generally, the preferred model has elevated the dextral rates of all faults in the Eastern California shear zone (Figure 16) above their target geologic rates, but only by an average of +0.6 datum standard deviations, so it has not exceeded 95%-confidence upper limits on other faults. The solution also incorporates high rates of distributed deformation (1–5 × 10^{-15}/s, Figure 6) to bring the net dextral rate up to the geodetic value. A third contribution comes from clockwise rotation of small crustal blocks in the northeast and east-central Mojave Desert, which is accommodated by left-lateral slip on E-trending faults separating these blocks. Another block which rotates clockwise is that containing Joshua Tree National Park, which lies between the Pinto Mountain and Blue Cut faults. These predicted clockwise rotation rates increase southward, from ~4°/m.y. just S of the Garlock fault, to ~10°/m.y. in the central eastern Mojave, and reach ~20°/m.y. in Joshua Tree National Park.

7.6. Walker Lane

Wesnousky [2005] presented a comprehensive review of active faulting and block rotation in the Walker Lane. Preferred model GCN2008088 supplements this with estimates of fault heave rates, as seen in Figure 15 and Figure 17.
The preferred model does not have any significant rate of slip on the Stateline dextral fault system [Guest et al. [2007]. Although this fault was assigned a target dextral rate of $2.4\pm9.1$ mm/a based on offset of $30\pm4$ km since $13.1\pm0.2$ Ma, its model rate is only $0.06$ mm/a. This is due to the lack of geodetic evidence for continuing strain in the region, and also to the lack of connecting structures on its SE end.

South of $37^\circ51'N$ (Boundary Peak, NV), the model has dextral shear shared between two widely-separated but Northward-converging fault systems. On the western system, dextral slip at $2.4$ mm/a on the Panamint Valley fault connects to dextral slip at $2.1$ mm/a on the Hunter Mountain-Saline Valley fault. Continuing northward, there is a gap before dextral slip is taken up by the White Mountains fault at mean rate $2.3$ mm/a; this gap is bridged by slip transfer to the nearby and parallel Owens Valley fault, which has a mean dextral component of $1.9$ mm/a.

The eastern dextral fault system includes (S to N) the Death Valley (So) train at $1.6$ mm/a, the Death Valley (Black Mountains frontal) train at $2.0$ mm/a (dextral component), the Death Valley (No) train at $2.6$ mm/a, and the Death Valley (N of Cucamongo) train at $1.4$ mm/a.

Both the western and the eastern systems have releasing bends (right steps) in the latitudes of Death Valley. This is the primary cause of the extensional fault-normal (D) rate component of $2.3$ mm/a predicted for the Death Valley (Black Mountains frontal) train. In this model, other normal or oblique-normal faults of the southern Walker Lane have relatively small extensional heave rates (e.g., Deep Springs fault $D = 0.6$ mm/a; So Sierra Nevada $D = 0.3$ mm/a), and do not form a connected extensional system.

As Wesnousky [2005] predicted, the central Walker Lane ($37^\circ51'\sim38^\circ25'N$) is occupied by the Excelsior-Coaldale block(s), bounded by ENE-trending faults with high rates of sinistral slip, which rotate clockwise at $\sim3^\circ$/m.y.. In the model these sinistral faults include (S to N): the connected Coaldale faults #1 & #2 at $2.5$ mm/a, connected sinistral faults #1302 and #1303 at $\sim1.5$ mm/a, and the sinistral faults of the southern Garfield Hills (#1304) at $3.5$ mm/a. (Fault names in this paragraph follow Haller et al. [2002].) Locally, the rotating block pulls away from the White Mountains block at $D = 2.3$ mm/a on the Boundary Peak detachment fault [dePolo, 1998]; isostatic rebound of the footwall probably explains the prominent height of this peak.

In the northern Walker Lane (Figure 17), there is another cycle of dextral faulting/block rotation/dextral faulting. At the border between Figures 16/17 ($38^\circ40'N$) the Gumdrop Hills fault ($3.1$ mm/a) and the Bettles Well-Petrified Springs fault ($1.6$ mm/a) are carrying most of the dextral slip [c.f. Wesnousky, 2005]. Then, in the greater Reno area ($39^\circ\sim40^\circN$, $120^\circ\sim119^\circW$) there is another set of 3 clockwise-rotating blocks bounded by 4 NE-trending sinistral faults ($1.1\sim2.8$ mm/a) including the Spanish Springs Peak fault. (Note that none of these faults has a dated offset feature to give a geologic rate, so NeoKinema has estimated these rates from kinematic compatibility.) Then, at $40^\circN$ (near the California border) dextral slip resumes, where it is divided between the Honey Lake fault ($1.2$ mm/a), Warm Springs Valley fault ($1.4$ mm/a), and Pyramid Lake fault ($2.3$ mm/a). This is not the end of the Walker Lane, as there are additional NW-trending faults in the lava beds of the Modoc Plateau which are presumably dextral or dextral-transtensional. However, as they were not catalogued in WGCEP Fault Models, their activity is treated here as distributed deformation, at rates up to $2\times10^{-15}$/s (Figure 6), extending up to the Oregon border.
7.7 Inland states

In the interior of the western U.S., geologic and geodetic data are less concentrated than they are in coastal states. Wesnousky et al. [2005] has collected geologic rates for many of the Basin & Range normal faults along one transect (Table 4), but elsewhere in the province Quaternary geologic rates are quite rare. Therefore, many Basin & Range normal faults in this model have target throw rates set to the generic $N = 0.183 \pm 0.343 \text{ mm/a}$ (see section 3.2), which (for assumed dip of 55°) implies a generic heave rate of $D = 0.128 \pm 0.24 \text{ mm/a}$. Other normal faults have rates based on minimum net throw (from scarp heights) since some time in the Oligocene or Miocene, but because of the throw uncertainty and great age of these offset features program Slippery attributes a comparable or even greater proportional uncertainty [Bird, 2007]. Also, the relatively low strain rates and fault slip rates in the interior mean that differences between velocities of adjacent geodetic benchmarks are typically less than the uncertainties in their velocities. These (relative) deficiencies in quantity and precision of inland data make the NeoKinema inverse problem “soft” or “easy” in the sense that many offset rates and geodetic velocities can be set to their prior/input values without causing any serious conflict with adjacent data. Consequently, the inland part of the map of long-term velocities (Figure 12) looks like a (smoothed) map of GPS velocities, while the maps of predicted fault heave rates (Figures 17-19) look very much like the target rates computed and tabulated in Table 1 of Bird [2007]. Due to length limits on this paper, these inland predicted fault heave rates are only presented graphically (and in attached digital supplements), without discussion.

However, one kind of artifact that appeared in these figures requires explanation. Preliminary versions of Figures 17-19 based on results from the preferred model GCN2008088 showed that certain Basin & Range “normal faults” (according to the input dataset) were predicted to be slipping as reverse faults. The worst region was east-central Nevada $(117° \sim 115° \text{W}, 39° \sim 41°)$, where 10 “normal faults” out of 195 had the wrong sense of throw, with the most negative rate at $-1.5 \text{ mm/a}$. This is very implausible from a dynamical point of view. Most likely these predictions are artifacts. They could result from geodetic velocities which are not “interseismic” as assumed, because there was some creep event or slow earthquake in the region [Davis et al., 2006; Wernicke et al., 2008], either during or just before the time window of geodetic observations. Alternatively, they could result from a “crowding” effect if the default normal throw rate is too high to apply in this region of many closely-spaced normal faults. Fortunately, such artifacts can be removed by an iterative process. In the first round of corrections, 22 inland “normal faults” with rapid reverse-slip predictions were removed from the input data for ad-hoc model GCN2008102, effectively locking these faults. During the new calculation in which these faults could no longer accommodate shortening, extensional rates on neighboring faults were reduced, and some became negative which had previously been positive. A second correction involving the deletion of 18 remaining wrong-way faults (even those with $N = -0.001 \text{ mm/a}$) gave ad-hoc model GCN2008103, which has very few artifacts in inland states. This iteratively corrected ad-hoc model is the basis for Figures 17-19.

8. Forecasting Seismicity

The techniques displayed here can contribute to forecasts of seismicity in 3 ways:

(1) Successful NeoKinema models provide better fault slip rates. Because geologic slip rates are unavailable, imprecise, or conflicting for many faults, committees of experts have often been assembled to choose rates (which I have referred to as “consensus composite rates”).
Program Slippery of Bird [2007] illustrates how computational statistics can be used to deal with conflicting or incomplete information about geologic offsets along any individual fault train. Program NeoKinema, run with input from Slippery, goes further by providing posterior/output estimates (predictions) of fault slip rates which also take into account geodetic velocities, stress directions, kinematic compatibility, and plate tectonics. Although slip rates predicted by one particular NeoKinema model do not come with uncertainties, a range of rates can be assembled from a suite of acceptable alternative models using different fault sets and Euler poles (etc.) as illustrated by attached file f_GCN_nko_ranges.txt (based on the 16 acceptable community models, 4 updated models, and 4 ad-hoc models of this paper). Expert panels would still be needed, but their roles could be modified to emphasize (a) collection and screening of data to be used in computations; (b) review of predictions for obvious artifacts, including (c) consideration of cases where NeoKinema predicts an unexpected sense of slip; and (d) consideration of paleoseismic studies as to whether particular faults creep or stick-slip.

The problems with extracting only improved slip-rates from NeoKinema modeling are that: (i) varying slip rates on one fault are typically correlated with varying slip rates on neighboring faults, so it is not appropriate to treat these refined slip-rate ranges as independent; (ii) long-term fault slip rates often vary along the fault trace, and this is not captured by using only the along-trace mean rate for seismicity prediction; and (iii) modeling in this paper has shown that as much as 1/3 of relative motion between some pairs of plates is accommodated by distributed permanent deformation off the mapped fault traces. Therefore, a superior approach considers that:

(2) Successful NeoKinema models are deformation models. Each computation provides estimates for the slip rate of each fault in each finite element (typically 15-30 km wide) as seen in Figures 13B~19 of this paper. These are “noisy” in two ways: they are discontinuous between elements, and sometimes implausibly large at fault terminations. However, it would be easy to apply smoothing if this were thought desirable. Also, each computation provides an estimate of the tensor of distributed permanent strain rate for each finite element. Again, discontinuities at element boundaries are artifacts, but these could easily be smoothed. (Any smoothing method should conserve seismic moment rate.) Using this fine-grained and detailed information from one finite-element model addresses all 3 concerns of the previous paragraph. (As examples, I attach 2 files as supplemental material with the predictions of the preferred model: h_GCN2008088.nko.txt, and e_GCN2008088.nko.txt.) The principal decision that has to be made is whether to consider distributed permanent deformation as a source of earthquakes; I argue that this is prudent.

Given a deformation model, there are still many controversial decisions which have to be made (or straddled) to get to a seismicity forecast. The report of 2007 WGCEP [2008] describes a complex logic-tree with branches expressing divergent opinions about fault segmentation, area/magnitude relations, characteristic earthquakes, periodic earthquakes, and ruptures outside the mapped fault traces. Incorporation of these many divergent models will always require a large team and very complex programs. This makes it difficult to update models quickly in response to new information, and it makes the modeling process laborious and expensive. In some cases, in may be desirable to consider a simpler two-step alternative:

(3) Program Long_Term_Seismicity can transform any preferred NeoKinema model into a map of estimated long-term seismicity. Then, the forecast can be made time-dependent by introducing the modulating effects of actual historical seismicity with an empirical statistical
model. The first step implies provisional acceptance of the SHIFT hypotheses reviewed in section 2.3 of this paper, and the calibration constants estimated by Bird & Kagan [2004]. The second step might involve treating some fixed fraction of the long-term seismicity map as the “background” or “immigrant” term in an epidemic-type earthquake sequence (ETES) model like that of Werner [2007], using maximum-likelihood methods to obtain the ETES parameters from the historic earthquake catalog, and then projecting the model forward in time. In this way, time-dependent processes such as aftershock sequences, earthquake clustering, and stress-shadowing could be added as perturbations to a steady process, while demonstrating at each step that the (single) model is statistically optimal and free of subjective elements.

To illustrate the first of these steps, I show in Figure 20 a calculation of long-term shallow seismicity based on preferred NeoKinema model GCN2008088 of this study, and computed with program Long_TERM_Seismicity (v.3, 2009.04.29). The forecast is also attached as supplemental material in digital form: LTSv3_GCN2008088_m5p663.grd.txt. The threshold is moment-magnitude 5.663 (scalar seismic moment $3.5 \times 10^{17}$ N m), above which the Global CMT catalog (like most seismic catalogs covering the western U.S.) is probably complete since 1977. Then, in Figure 21 I superpose actual shallow seismicity from 1977.01.01-2008.11.30 in the same region from the Global CMT catalog. There is absolutely no circularity in this comparison because historical seismicity played no part in the NeoKinema modeling, and because the Gorda-California-Nevada orogen was excluded from the spatial domain used by Bird & Kagan [2004] to estimate the seismicity constants of different kinds of plate boundaries.

In a longitude/latitude “trapezoid” (128~104°W, 30~49°N) surrounding the NeoKinema model, the forecast long-term seismicity rate of $m>5.663$ shallow earthquakes is 3.54/year. The actual record in 1977.01.01-2008.11.30 was 71 earthquakes, or 2.22/year. This suggests that the western U.S. has been 37% below its long-term seismicity rate, and should be expected to have more shallow earthquakes in the future. The map pattern in Figure 21 shows several prominent seismic gaps: First, the Cascadia subduction zone (as opposed to the adjacent Gorda orogen) in the smaller trapezoid (128~122°W, 42~49°N) produced only 9 $m>5.663$ shallow earthquakes in this period for a rate of 0.28/year, although its long-term average rate is predicted to be 0.95/year. In the remainder of the large trapezoid, the deficit was less dramatic: 1.94/year actual versus 2.57/year expected (76% of expectation). Figure 21 shows that much of this remaining deficit occurred along the North Coast, Big Bend, and northern Gulf of California portions of the Pacific-North America boundary.

We know from historic great earthquakes in 1700 AD (Cascadia subduction zone) and 1857 and 1906 AD (San Andreas fault) that the relative quiescence in at least 3 of these regions is temporary. Whenever the next great earthquake occurs on either fault, it is likely to be associated with clusters of $m>5.663$ aftershocks and more distant triggered seismicity which will make up the deficit, and likely even exceed the long-term rate for several decades.

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Figure 1. Finite element grid GCN8p9.feg used in this study. Most of the grid is composed of quasi-equilateral spherical triangles, with sides of either 30 km (coarse regions) or 15 km (fine regions). Ribbons of smaller elements, with width approximately 4 km, have been inserted along most fast-slipping faults to better approximate the expected velocity discontinuities. There are 6452 nodes and 12627 elements.
Figure 2. Traces of 1472 active and potentially-active faults included in these models. Traces are colored according to prior expectations of their predominant sense(s) of slip. Faults with oblique slip have a green or brown trace to indicate dextral or sinistral component, plus dip-ticks of a different color and shape to indicate the primary mode of dip-slip. Offset type D is used for both low-angle detachment faults and magmatic spreading centers.
Figure 3. GPS benchmarks, interseismic velocities, and 95%-confidence ellipses used in modeling. As described in text, California velocities are from a 2006 solution by Shen and others for WGCEP; velocities outside California are selected from the PBO solution of September 2007. All velocities are in the stable North America reference frame. Guadalupe Island is just visible at the southern margin of the map.
Figure 4. Data on the azimuth of the most-compressive horizontal principal stress from the World Stress Map (A), and directions interpolated by NeoKinema (B) using the clustered-data algorithm of Bird & Li [1996].
Figure 5. Neotectonic Euler poles for relative rotation of North America (NA) with respect to Pacific (PA). The error ellipses shown are standard errors, so 95%-confidence ranges have twice the diameters shown, and typically overlap. For each pole, a label indicates the implied NA-PA relative velocity at Parkfield, CA (assuming that stable NA and PA lithosphere extend up to the San Andreas fault at that point, whereas actually they do not). Poles within the dashed rectangle were used in NeoKinema modeling; others are shown for historical interest. The Gulf_GPS pole was not explicitly stated by Antonelis et al. [1999] but was computed by the author from their velocity vectors.
Figure 6. Common logarithms of distributed permanent strain-rates (excluding strain-rates due to slip on modeled faults) in the preferred model GCN2008088. See equation (5) for definition of scalar measure $\dot{e}$. 
Figure 7. Three misfit measures ($N_{\text{geodetic}}^2$, $N_{\text{potency}}^2$, $N_{\text{stress}}^2$) are contoured in a 2-D parameter space with axes of $\log_2(L_0)$ and $\log_2(A_0)$. Contour interval 0.2, with heavier lines at value 2.0, and colored shading to show regions of unacceptable misfit (any $N_2 > 2$). Acceptable models are shown by rectangles and octagon, while unacceptable models are shown by triangles. All computations used prior/input $\mu = 5 \times 10^{-16}$, Fault Model 2.1, NUVEL-1A pole, and block-diagonal approximation of the geodetic covariance.
Figure 8. Posterior/output values of RMS distributed permanent strain-rate ($\mu^* = \text{RMS}(\dot{e})$) shown with contours in the same 2-D parameter space as Figure 7. All inputs as in Figure 7. Acceptable models (with all misfit measures < 2) are shown by rectangles and octagon, while unacceptable models are shown by triangles.
Figure 9. Posterior/output values of RMS distributed deformation rate ($\mu^*$, in A) and 3 misfit measures ($N_2^\text{geodetic}$, $N_2^\text{potency}$, $N_2^\text{stress}$, in B) plotted as functions of input parameter $\mu$, with fixed weights ($L_0=4\times10^4$ m, $A_0=4\times10^8$ m$^2$), and other inputs as in Figure 7. Note that output $\mu^*$ is relatively insensitive to input $\mu$, and that this problem has a natural minimum $\mu^*$ of $5\times10^{-16}$/s.
Figure 10. Three misfit measures ($N_2^{\text{geodetic}}$, $N_2^{\text{potency}}$, $N_2^{\text{stress}}$) are contoured in a 2-D parameter space with axes of $\log_2(L_0)$ and $\log_2(A_0)$. All conventions as in Figure 7. The differences here are that the full covariance matrix of California GPS velocities is used, the NA-PA Euler pole is the Guadalupe pole, and southern California fault traces are from WGCEP Fault Model 2.2.
Figure 11. A pseudo-prospective test of the ability of the set of 16 successful “community” models to predict “new” long-term geologic offset rates which were not used in their computation. Data sources in Table 4. Large discrepancies are discussed individually in text Section 7.
Figure 12. Long-term velocity field of the preferred model GCN2008088. Note that effects of transient elastic strain accumulation about the Cascadia trench and San Andreas fault system (and all other faults) have been removed. Brightness contour interval 1 mm/a; jagged contours are caused by velocity discontinuities across faults. For legibility, velocity vectors are shown at only 1/9 of nodes. Velocity reference frame is stable eastern North America.
Figure 13. Fault heave rates from preferred model GCN2008088 in the Washington-Oregon region, displayed in two formats: (A) The trace-averaged heave rate is plotted at every point along the trace, giving ribbons of uniform width. (Oblique slip is represented by two ribbons of different colors plotted along the same trace.) (B) Individual per-element heave rates are plotted, without enforcing continuity along trace. This “noisy” plot has the potential advantage of displaying predicted variations in offset rate along each trace. However, it also displays probable artifacts, such as implausible high rates in elements where faults terminate without any fault junction.
Figure 14. Fault heave rates predicted by NeoKinema in the region of the Mendocino triple junction. (A) Preferred model GCN2008088, in which the Mendocino fault is allowed to slip obliquely and absorbs 10 mm/a of N-S shortening by underthrusting Gorda crust under Pacific. (B) Ad-hoc model GCN2008101 in which the Mendocino fault is vertical, and shortening takes place by distributed deformation, faster sinistral faulting within the Gorda crust, and its accelerated subduction at the south end of the Cascadia trench.
Figure 15. Fault heave rates from preferred model GCN2008088 in the San Francisco Bay, central California Coast Ranges, and central and southern Walker Lane regions.
Figure 16. Fault heave rates from preferred model GCN2008088 in southern California.
Figure 17. Fault heave rates from model GCN2008103 in the northern Walker Lane and northern Nevada. In the Walker Lane, fault traces have been overlain on the heave-rate ribbons of 4 NE-trending sinistral faults. Elsewhere, fault traces are not overlain because they would obscure small heave-rates.
Figure 18. Fault heave rates from model GCN2008103 in the northern Rocky Mountains and northeastern Basin & Range province. While few faults are mapped within the Snake River Plain (shaded), it is moving [Payne et al., 2008] and deforming by other means [Parsons et al., 1998].
Figure 19. Fault heave rates from model GCN2008103 in the regions surrounding the Colorado Plateau.
Figure 20. Common logarithm of long-term shallow seismicity (epicenters per square meter per second, including aftershocks) for threshold magnitude 5.663 (moment $3.5 \times 10^{17}$ N m), computed by Long_Term_Seismicity (v.3) from preferred NeoKinema model GCN2008088. Seismicity around the margins, outside the NeoKinema model domain, is based on relative plate motions from model PB2002 of Bird [2003] and intraplate strain rates from dynamic Shells model Earth5-049 of Bird et al. [2008]. Rates in central Montana and eastern Wyoming are too high, for reasons explained in that paper. The spatial integral of these forecast rates is 113 earthquakes per 31.92 years in the depth range 0–70 km. (To convert seismicity from earthquakes/m²/s to
earthquakes/km$^2$/year, add 13.5 to the values along the logarithmic scale. To convert to earthquakes/(100 km)$^2$/century, add 19.5.)
Figure 21. Colored background shows long-term forecast, exactly as in Fig. 20. For retrospective comparison, the Harvard CMT catalog shows 71 events (with focal mechanisms on lower focal hemispheres) of $m>5.663$ at 0~70 km depth in the 31.92-year interval 1977.01.01~2008.11.30. This figure illustrates why the instrumental record of seismicity is very inadequate for estimating maps of long-term seismicity.


Table 1. Comparison of modeling methods, input data counts, misfit measures, and numbers of predictions

<table>
<thead>
<tr>
<th>Kinematic model of neotectonics of/within/including the western U.S.</th>
<th>Model type (a)</th>
<th>Area, 10^{12} m^2</th>
<th>Elements/cells/blocks</th>
<th>RMS resolution, km</th>
<th>Input geologic offset rates (b)</th>
<th>Input geodetic benchmarks</th>
<th>Input stress azimuths</th>
<th>( \sum \chi^2/n ) of best model</th>
<th>Predicted fault offset rates (c)</th>
<th>Predicted off-fault permanent strain-rate tensors</th>
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<tr>
<td>Saucier &amp; Humphreys [1993]</td>
<td>F-E</td>
<td>0.36</td>
<td>400</td>
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<td>10</td>
<td>0</td>
<td>?</td>
<td>33</td>
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<td>5</td>
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</table>

(a) Block = purely-elastic microplates; Soft-block = microplates with 3 DOF each for internal permanent strain-rate; F-E = finite-element grid; Spline = velocity derived from Euler pole, and Euler pole components interpolated laterally by splines on a deformed quadrilateral grid; VE = analytical viscoelastic model with faults.

(b) Counted as number of fault trains with at least one dated offset feature supporting the target offset rate for that train, not as total number of offset features. Fault trains with generic/default target offset rates are not counted.

(c) From 1 to 3 offset-rate components per fault train; in this study only 1 or 2 per train. In cases of Block and Soft-block models, components may include nonphysical fault-orthogonal components on vertical faults.

(d) Consensus composite fault slip rates selected by a committee of experts are influenced by seismicity, paleoseismicity, geodesy, plate tectonics, and geometric compatibility as well as by dated offset geologic features (if any). Thus, many do not meet the criteria for counting in note (b).

(e) After “iterative elimination of outliers” [Meade & Hager, 2005, paragraph 26].
Table 2. Alternative neotectonic Euler poles for NA-PA relative rotation

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<tr>
<th>Name</th>
<th>Reference(s)</th>
<th>N. lat.(deg.)</th>
<th>E. lon(deg.)</th>
<th>Rate(deg./m.y.)</th>
<th>Velocity (mm/a)</th>
<th>Azimuth (deg.)</th>
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<tr>
<td>NUVEL-1A</td>
<td>DeMets et al. [1990; 1994]</td>
<td>48.709</td>
<td>-78.167</td>
<td>0.7486</td>
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<td>Gulf_GPS*</td>
<td>based on Antonelis et al. [1999]</td>
<td>51.7</td>
<td>-81.1</td>
<td>0.746</td>
<td>43.9</td>
<td>137.9</td>
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<td>REVEL-2000*</td>
<td>Sella et al. [2002]</td>
<td>50.38</td>
<td>-72.11</td>
<td>0.755</td>
<td>50.9</td>
<td>141.8</td>
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<td>ITRF2000</td>
<td>Altamimi et al. [2002]</td>
<td>50.488</td>
<td>-75.134</td>
<td>0.755</td>
<td>48.7</td>
<td>141.3</td>
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<td>PA_GPS</td>
<td>Beavan et al. [2002]</td>
<td>50.26</td>
<td>-75.04</td>
<td>0.773</td>
<td>49.9</td>
<td>141.7</td>
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<td>Guadalupe</td>
<td>Gonzalez-Garcia et al. [2003]</td>
<td>49.89</td>
<td>-77.01</td>
<td>0.766</td>
<td>47.8</td>
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<td>ITRF2005</td>
<td>Altamimi et al. [2007]</td>
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<td>0.773</td>
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<td>Plate_frame</td>
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<td>51.16</td>
<td>-73.83</td>
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*illustrated in Figure 5, but not used in modeling.
Table 3. Computed models and their misfit measures

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<th>Model</th>
<th>NA-PA pole</th>
<th>CA Fault</th>
<th>GPS covariance?</th>
<th>$L_0$, $A_0$, $\mu$, $\mu^*$, $N_2^{\text{geodetic}}$, $N_2^{\text{potency}}$, $N_2^{\text{stress}}$</th>
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<th>1.6E+09</th>
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<td>2.685</td>
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<th>Plate Frame</th>
<th>Resolution</th>
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<th>T0 (s)</th>
<th>T0 (s)</th>
<th>T1 (s)</th>
<th>T1 (s)</th>
<th>T2 (s)</th>
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Table 4. Fault offset rates predicted by acceptable community models compared to new geologic offset rates

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<th>Lat., °</th>
<th>Fault, State</th>
<th>Neo-Community Model</th>
<th>New geologic offset rate (Slippery, 95% CI)</th>
<th>Updated Models: Discrepancy Prediction: Discrep.</th>
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<td>F2337</td>
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<td>Blue Lake thrust fault, CA</td>
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<td>1–1.45 McCrory [2000] (site 7)</td>
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<td>East Humboldt Range normal fault, NV</td>
<td>F0512</td>
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<td>0.294 Wesnousky &amp; Willoughby [2003]</td>
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<td>40.82</td>
<td>Independence Valley normal fault, NV</td>
<td>F2138</td>
<td>N</td>
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<td>40.76</td>
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<td>40.75</td>
<td>Wasatch normal fault, UT</td>
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<td>D</td>
<td>0.616–0.626 Armstrong et al. [2004]</td>
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<td>40.64</td>
<td>Buena Vista (Beachfront) normal fault, NV</td>
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<td>Max Depth (m)</td>
<td>Step Length (m)</td>
<td>Slip Rate (cm/yr)</td>
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<td>0.15–0.151</td>
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<td>Dry Hills(?) normal fault, NV</td>
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<td>Stansbury normal fault, UT</td>
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<td>0.299–0.3</td>
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<td>0.901–1.42</td>
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<td>F2225</td>
<td>1.2</td>
<td>1.22–1.26</td>
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<td>Monte Cristo Valley dextral fault, NV</td>
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